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Introduction to neutron reflection

Adrian Rennie



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Outline

Inteference of waves

Refractive index

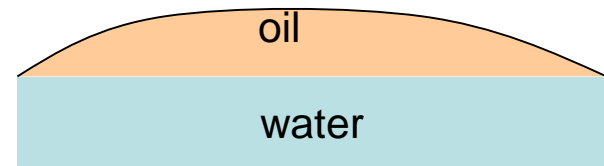
Critical angle, total reflection



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Reflection

Light

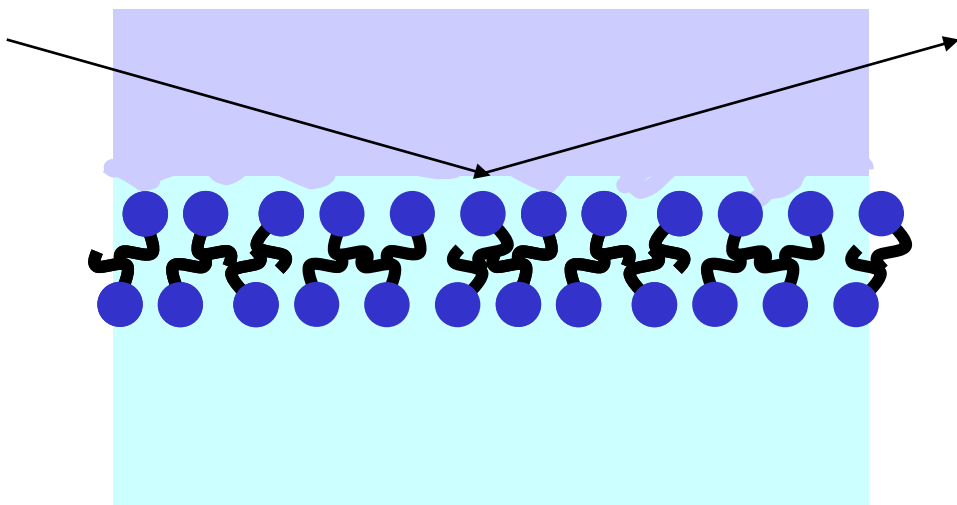




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Reflection

Light

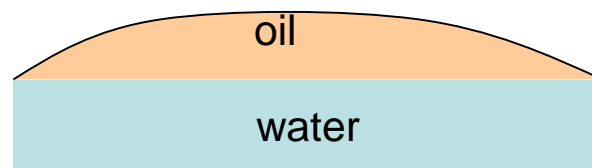


Why neutrons?

Contrast: light elements, isotopes

Penetrate

Magnetism





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Reflection and Refraction: Snell's Law

For specular reflection:

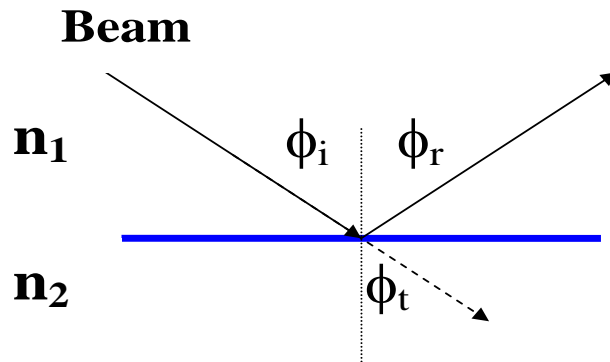
$$\phi_i = \phi_r$$

Transmitted beam is
refracted:

$$n_2 \sin \phi_t = n_1 \sin \phi_i$$

n is refractive index

Optical Notation





Reflection and Refraction: Snell's Law

For specular reflection:

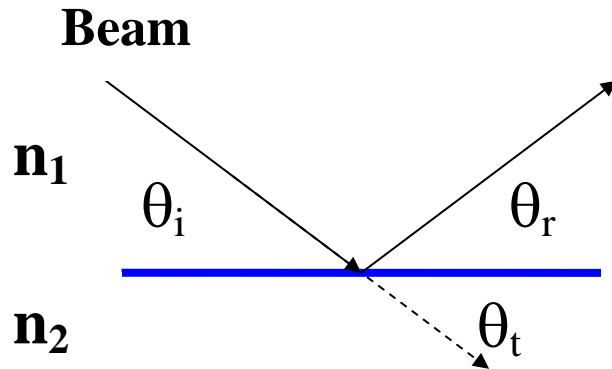
$$\theta_i = \theta_r$$

Transmitted beam is refracted:

$$n_2 \cos \theta_t = n_1 \cos \theta_i$$

n is refractive index

Neutron Reflection
Notation



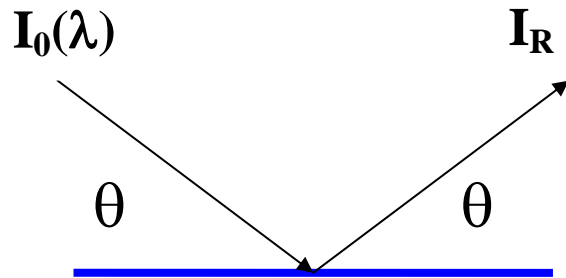
$$\theta = 90^\circ - \varphi$$



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Reflection – measured quantities

Reflection



Reflected beam
deflected: 2θ

Reflectivity

$$R(Q) = I_R / I_0(\lambda)$$

Momentum transfer

$$Q = (4\pi/\lambda) \sin \theta$$



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Demonstration Calculations

www.ncnr.nist.gov/instruments/magik/calculators/reflectivity-calculator.html

www.ncnr.nist.gov/instruments/magik/calculators/magnetic-reflectivity-calculator.html



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Critical Angle and Below (critical wavelength and above)

Density difference between two bulk phases determines the critical momentum transfer/angle, Q_c or θ_c

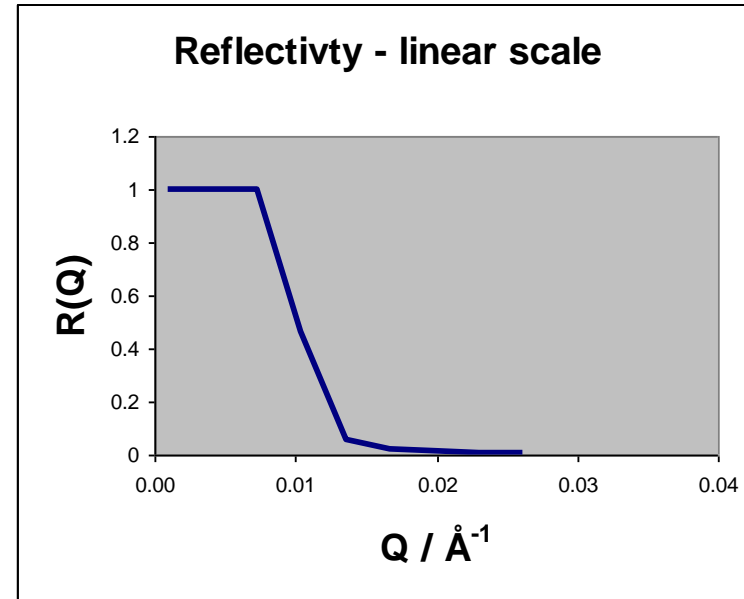
Any variation in intensity below critical angle is probably telling you about the experiment rather than the interface

$R(Q) = 1$ for $\theta < \theta_c$ is often used as a calibrant

$R(Q) \sim 1/Q^4$ for sharp interface

Total reflection below critical angle θ

$$\cos \theta = n_2/n_1$$





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Calculating Refractive Index

Neutrons

$$n = 1 - (\lambda^2 \Sigma_i b_i / V / 2\pi)$$

λ is the wavelength

$\Sigma_i b_i$ is the sum of scattering lengths in volume V

b is known for most stable nuclei

$$\rho = \Sigma_i b_i / V$$



Scattering Lengths of Nuclei

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Nucleus	Scattering Length / fm
^1H	-3.741
^2H (or D)	6.675
C	6.648
O	5.805
Si	4.151
Cl	9.579

Source: H. Rauch & W. Waschowski



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Properties of Common Materials

Material	Scatt. Length Density / 10^{-6} \AA^{-2}	Refractive index at 10 \AA
H ₂ O	-0.56	1.000009
D ₂ O	6.35	0.999899
Si	2.07	0.999967
Air	0	1.000000
Polystyrene	1.4	0.999971



Contrast in a Thin Film

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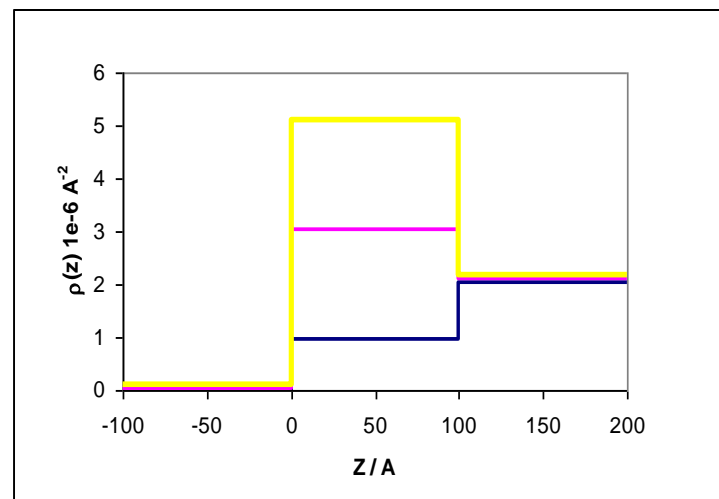
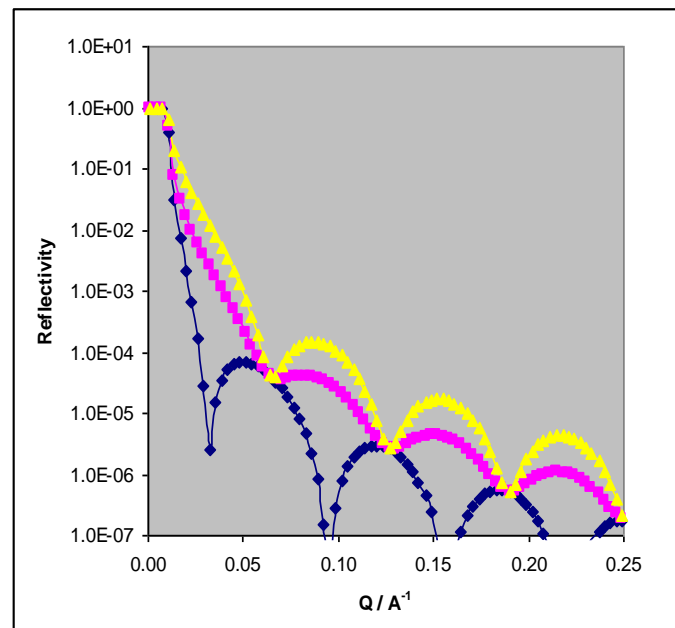
Calculation for Neutrons

100 Å layer with $\rho=1, 3 \text{ \& } 5 \times 10^{-6} \text{ \AA}^{-2}$
on Si ($\rho=2.07 \times 10^{-6} \text{ \AA}^{-2}$)

Increasing contrast changes visibility of fringes

Phase change makes large difference

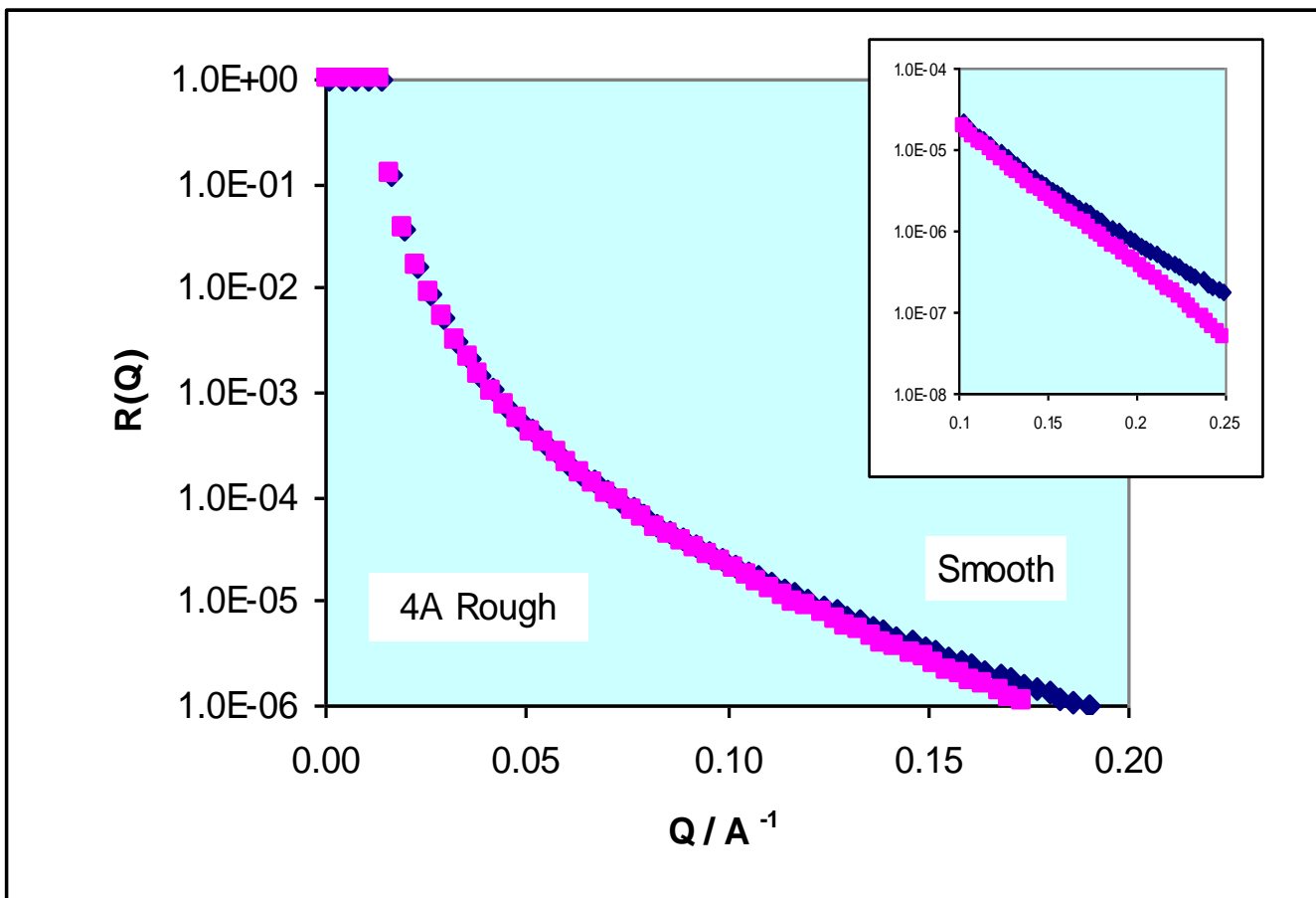
Fringes (Kiessig fringes) – spacing indicates film thickness for a single layer.





Roughness

Reflectivity from rough surfaces is decreased.





Intensity of Reflected Signal

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Waves interfere constructively for

$$2 d \sin \theta = \lambda, 2\lambda, 3\lambda \dots \text{ (Bragg's law)}$$

Measured reflectivity will depend on angle and wavelength.

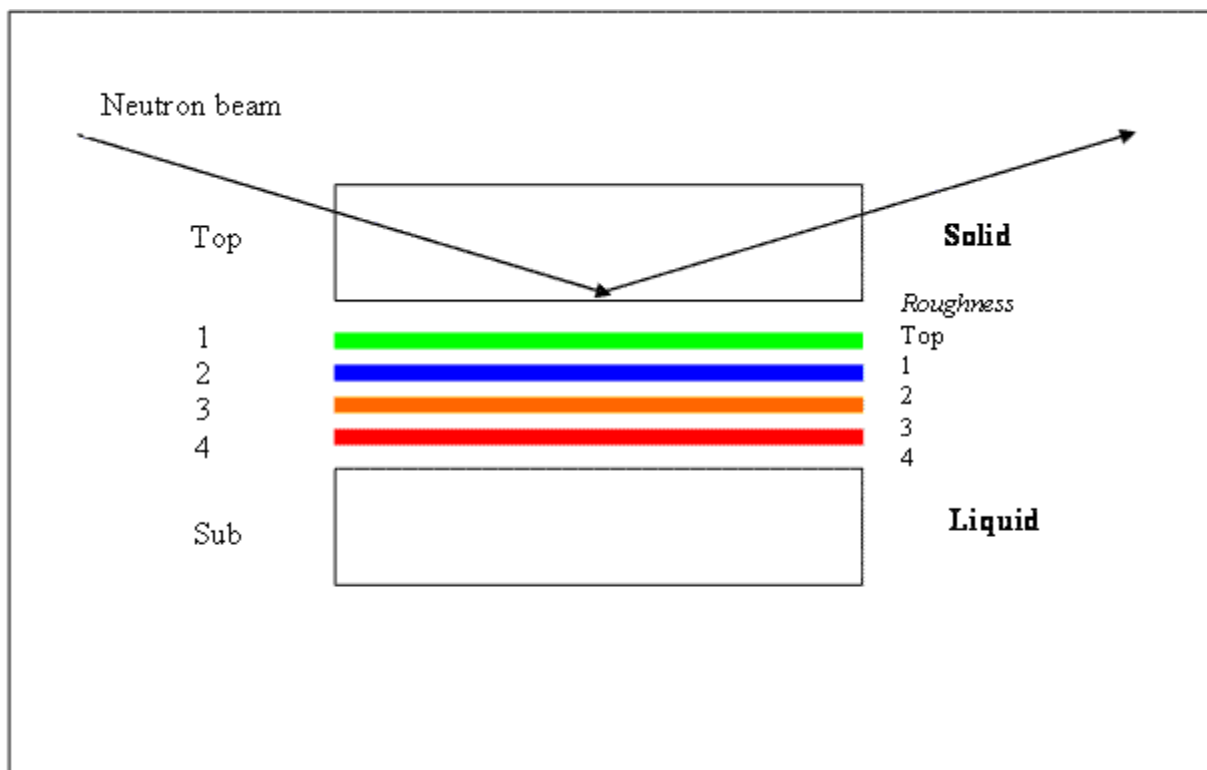
Total reflection for angles less than critical angle, $\theta_c = \arccos(n_1/n_2)$



Useful Physical Ideas

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Models for complex interfaces can be constructed from multiple thin layers of different refractive index, n or scattering length density, ρ .





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Useful Physical Ideas

Isotopes (e.g. D/H substitution) can be used to label particular species or alter contrast

Neutrons have spin – effectively a field dependent contribution to scattering length



Abeles Optical Matrix Method

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$$r_j = \begin{bmatrix} e^{i\beta_{j-1}} & r_{j-1}e^{i\beta_{j-1}} \\ r_{j-1}e^{-i\beta_{j-1}} & e^{-i\beta_{j-1}} \end{bmatrix}$$

$$\beta_j = (2\pi/\lambda)n_j d_j \sin \theta_j$$

$$p_j = n_j \sin \theta_j$$

$$r_j = (p_{j-1} - p_j)/(p_{j-1} + p_j)$$

$$M_R = [M_1][M_2]\dots[M_{n-1}]$$

$$R(Q) = M_{21}M_{21}^* / M_{11}M_{11}^*$$



Magnetic Contrast

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$$b_m = \mu_0 e^2 S \gamma / 4\pi m_e$$

e , electronic charge

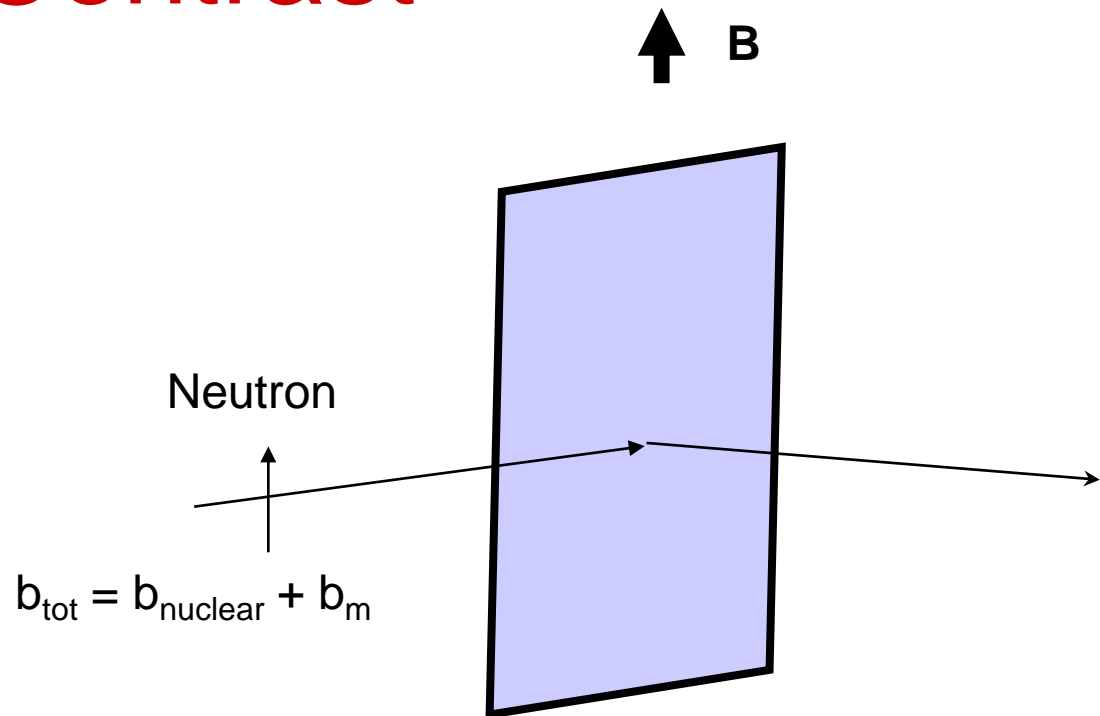
m_e , electron mass

S , spin

μ_0 , Permeability of free space

γ , gyromagnetic ratio, 1.913

$$b_{\text{tot}} = b_{\text{nuclear}} \pm b_m$$





Magnetic Contrast

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$$b_m = \mu_0 e^2 S \gamma / 4\pi m_e$$

e , electronic charge

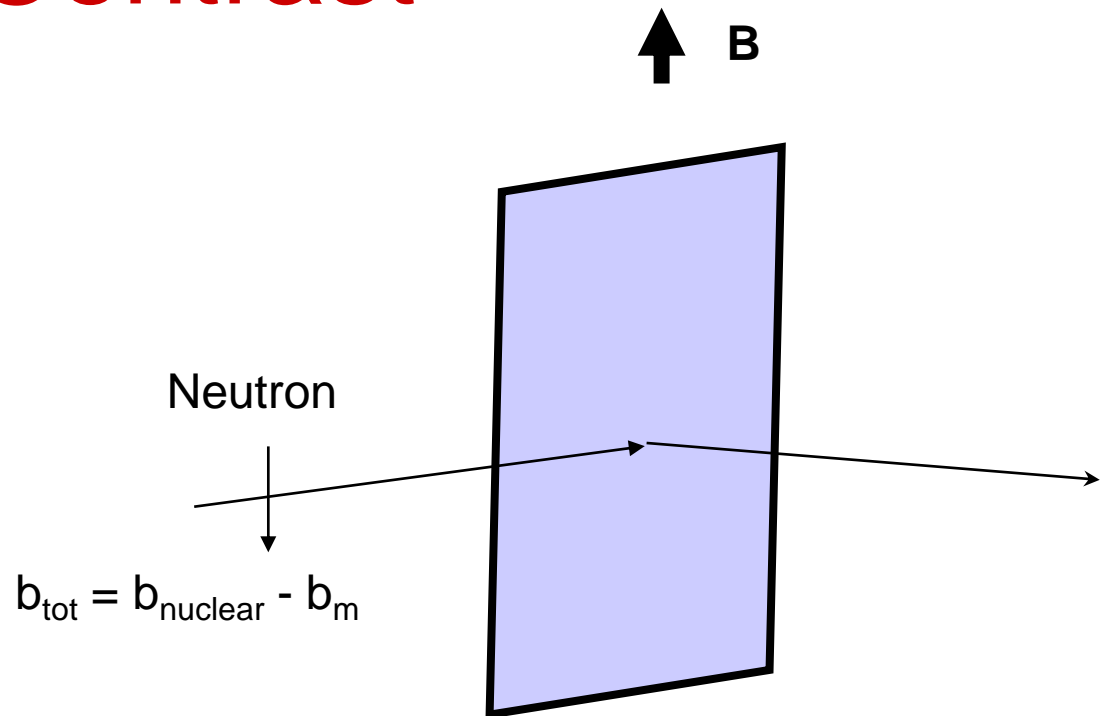
m_e , electron mass

S , spin

μ_0 , Permeability of free space

γ , gyromagnetic ratio, 1.913

$$b_{\text{tot}} = b_{\text{nuclear}} \pm b_m$$





Scattering and Reflection

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$\rho(Q)$ is Fourier
transform of the
scattering length
density distribution
normal to the interface,
 $\rho(z)$

$$R(Q) = \frac{16\pi^2}{Q^2} |\rho(Q)|^2$$

$$\rho(Q) = \int_{-\infty}^{\infty} \rho(z) e^{-iQz} dz$$

For sharp interface:

$$R(Q) \sim 1/Q^4$$



Partial Structure Factors

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Interface consists of distinct components: 1, 2, 3

$$R(Q) = \frac{16\pi^2}{Q^2} \left| \int \rho(z) e^{iQz} dz \right|^2$$

$$\rho(z) = b_1 n_1(z) + b_2 n_2(z) + b_3 n_3(z)$$

$$R(Q) = \frac{16\pi^2}{Q^2} (b_1^2 h_{11} + 2b_1 b_2 h_{12} + b_2^2 h_{22} + 2b_2 b_3 h_{23} + b_3^2 h_{33} + 2b_3 b_1 h_{31})$$

h_{ij} are transforms of $n_i n_j$ – pair correlation functions



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Practical Aspects of Neutron Reflection

How to Collect Data

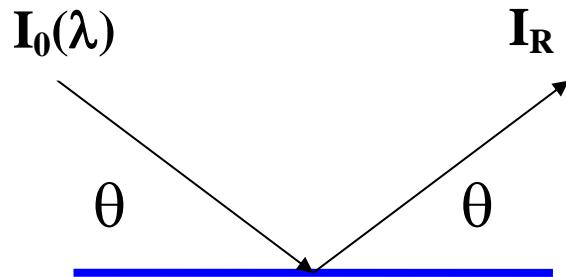
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Reflection – measured quantities

Reflection



Reflected beam deflected:

$$2\theta$$

Reflectivity

$$R(\theta, \lambda) = I_R / I_0(\lambda)$$

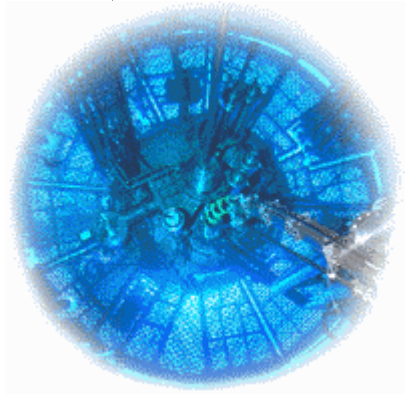
Momentum transfer

$$Q = (4\pi/\lambda) \sin \theta$$



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Best Sources of Neutrons



ILL reactor continuous

Thermal Flux $1.5 \times 10^{15} \text{ n cm}^{-2} \text{ s}^{-1}$



SNS, ORNL

60 Hz, 300 μs

$5 \times 10^{17} \text{ n cm}^{-2} \text{ s}^{-1}$ (Peak)



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Neutrons: Speed & Wavelength

Velocity, v , from de Broglie relation

$$v \lambda = 3956 \text{ m s}^{-1} \text{ \AA}$$

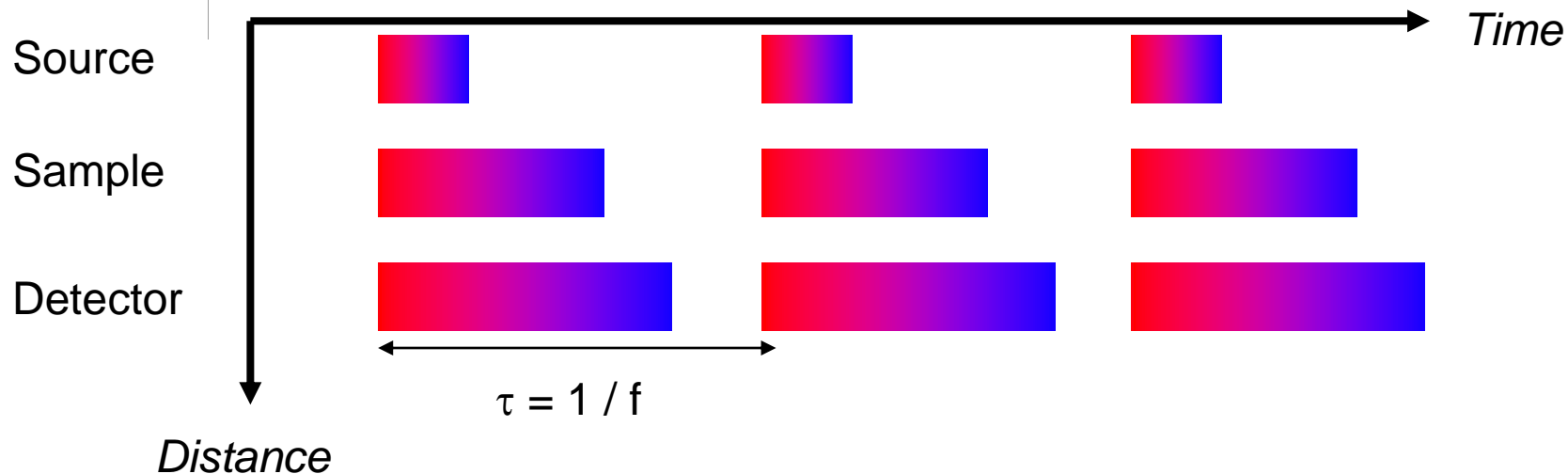
i.e. 10 \AA has 400 m s^{-1}

Gravity is significant, separate wavelengths mechanically



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Using a Pulsed Source



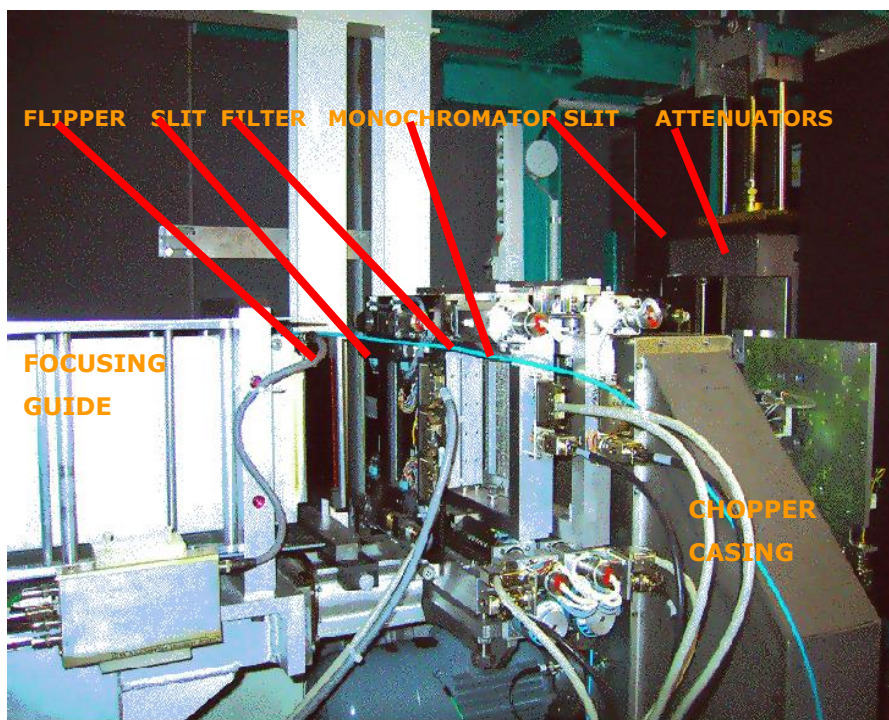
Detection time (after source pulse) gives wavelength

Choppers can select a wavelength



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D17 Reflectometer





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Practical Issues

Reflectivity drops quickly with increasing Q (or angle). Signal is easily 'lost' in background.

To observe fringes it will be necessary to measure over an appropriate range of Q and to have sufficient resolution (ΔQ small enough).



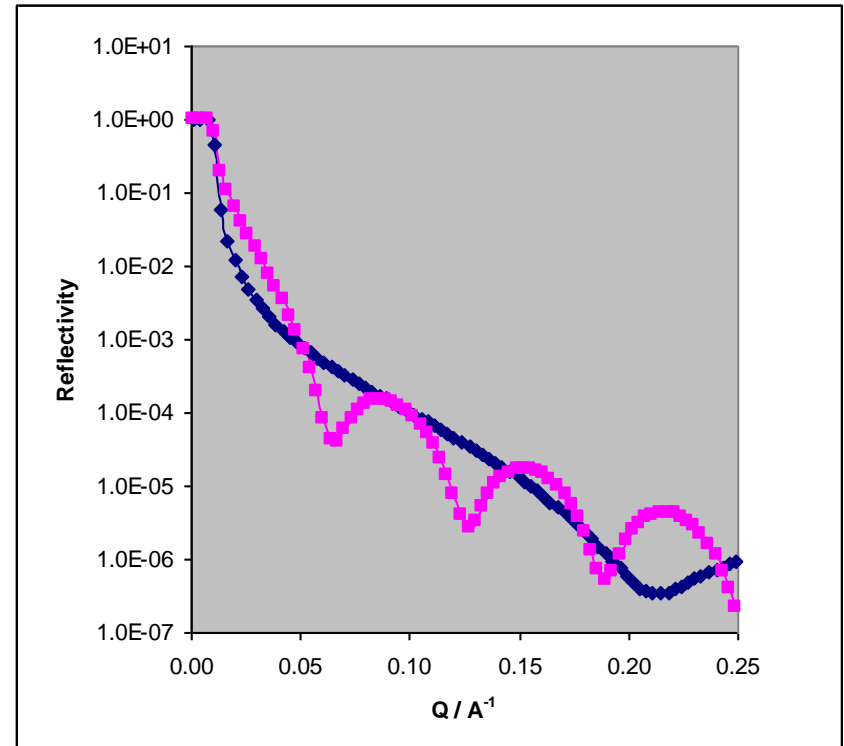
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Reflection from a Thin Film

Model calculation on
smooth surface.

Fringe spacing depends
on thickness

Fringe spacing $\sim 2\pi/d$



Model layer with $\rho = 5 \times 10^{-6} \text{ \AA}^2$ on Si
($2.07 \times 10^{-6} \text{ \AA}^{-2}$) Blue 30 \AA , Pink 100 \AA .
No roughness.



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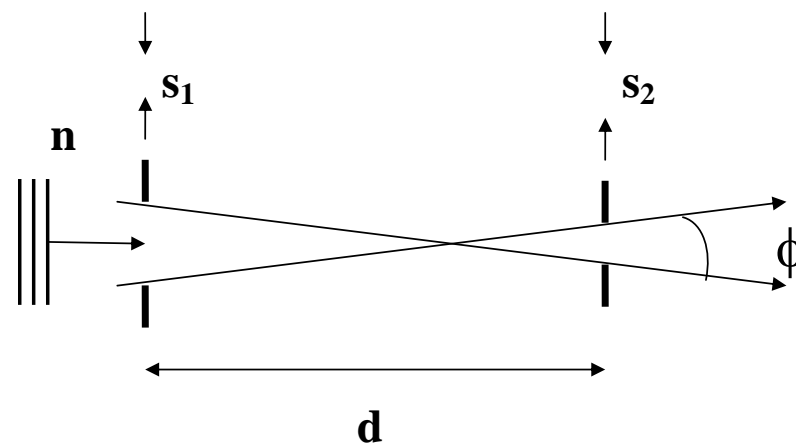
Resolution in Q

$$Q = (4\pi/\lambda) \sin \theta$$

Depends on $\Delta\lambda$ and $\Delta\theta$

Angle resolution, $\Delta\theta$, depends on collimation (slits)

Wavelength resolution depends on monochromator or time resolution in measuring neutron pulse



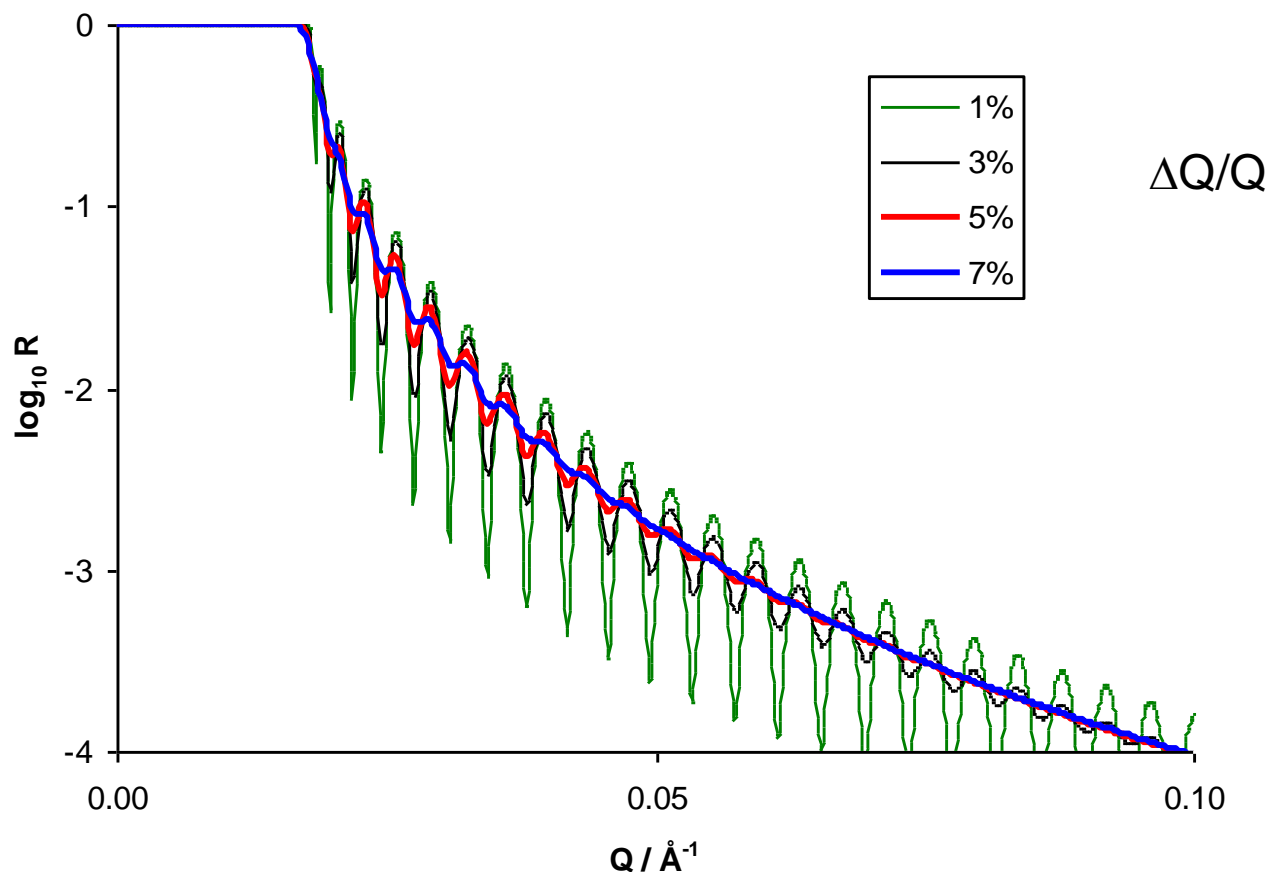
Higher Resolution = Lower Flux

$$(\Delta Q/Q)^2 = (\Delta\lambda/\lambda)^2 + (\Delta\theta/\theta)^2$$



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Effects of Resolution



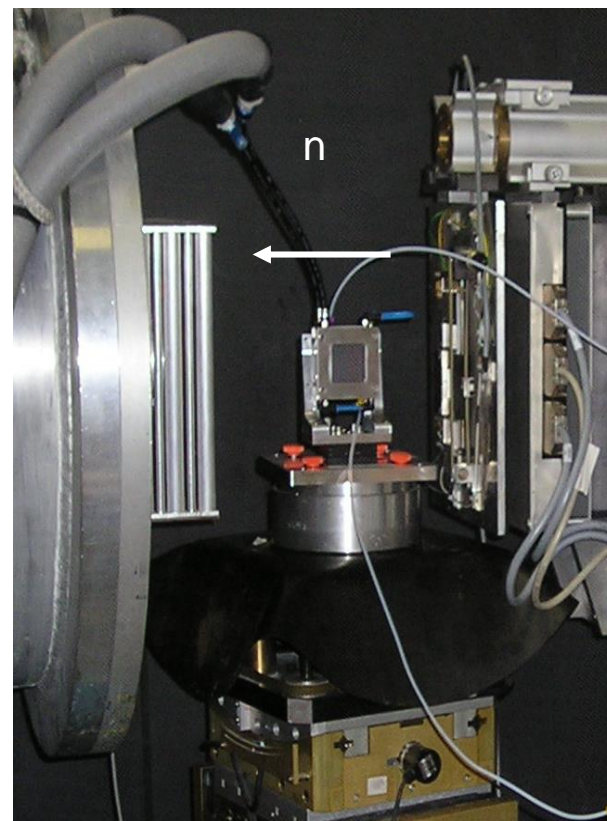
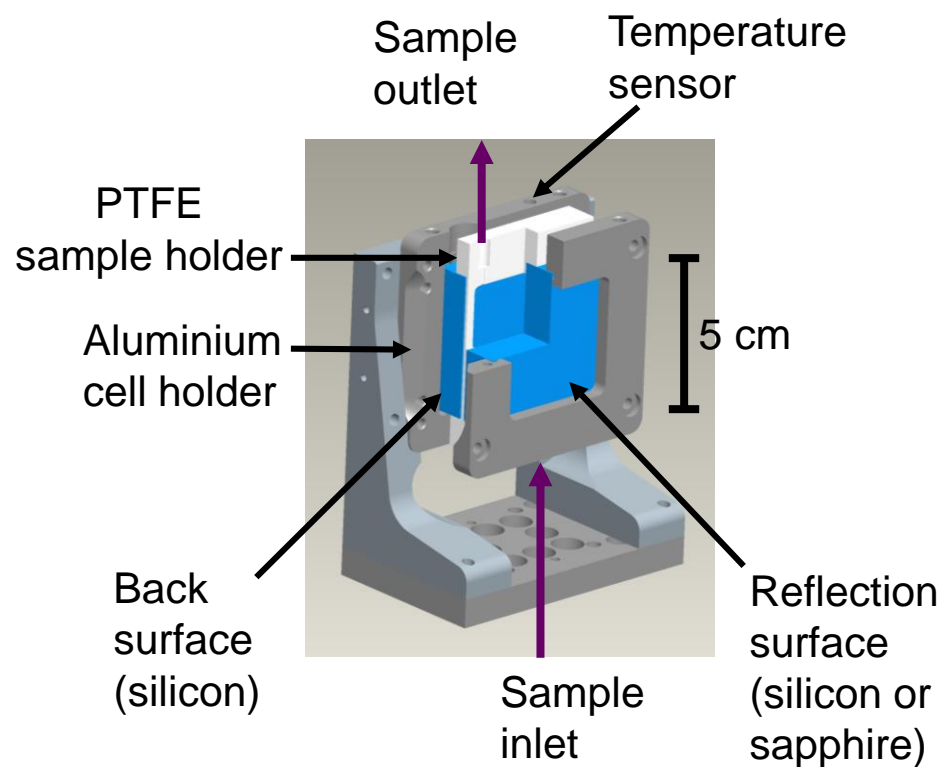
Silicon substrate: film thickness 1500 \AA (150 nm)
scattering length density $6.3 \times 10^{-6} \text{\AA}^{-2}$



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Sample Holder

D17 reflectometer
ILL, France





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Alignment

Rotation table must have centre on beam axis

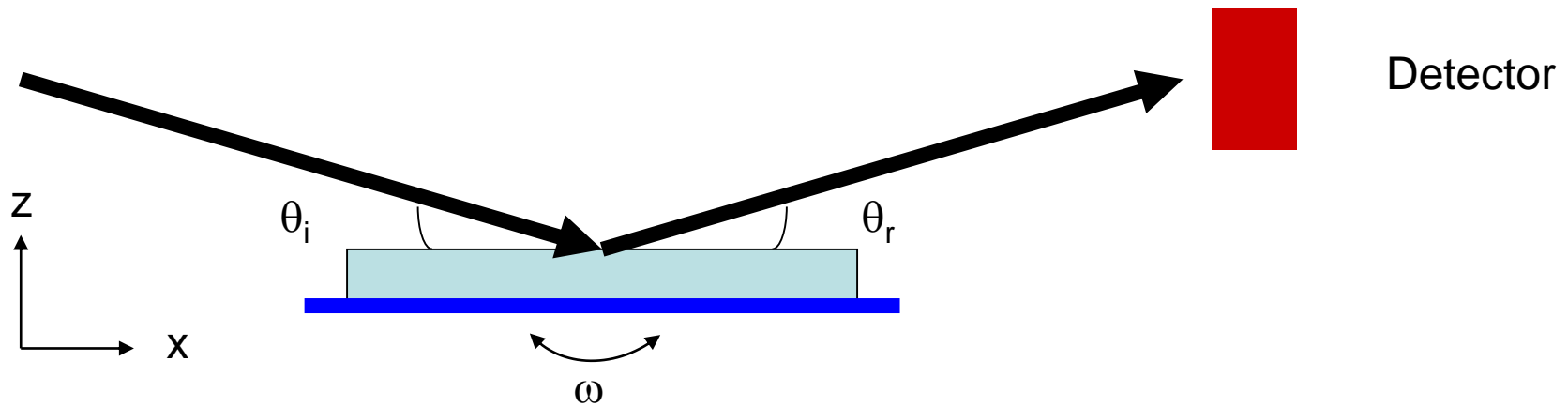
Sample must be centred on rotation (half obscure the direct beam) – eucentric mount

Determine θ from the position of beam on a detector



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Aligning a Sample



Design mount with surface at centre of rotation of ω . Eucentric mount.

Put centre of surface on the line through axis of rotation (x direction)

The rotation ω stage must be centred on the incident beam.

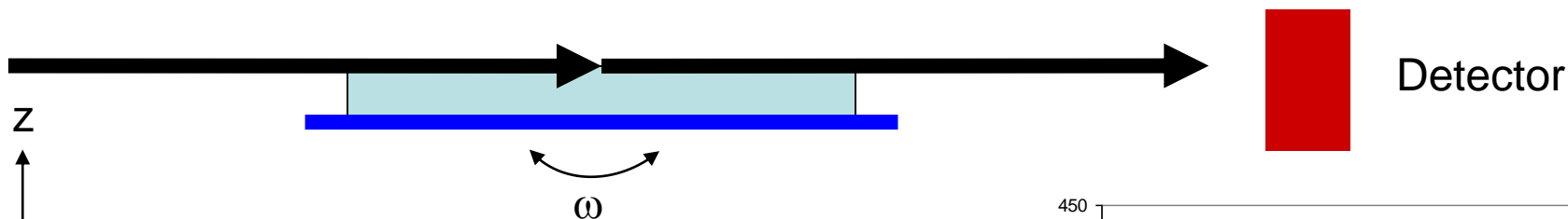


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Aligning a Sample

Set sample and detector to nominal zero

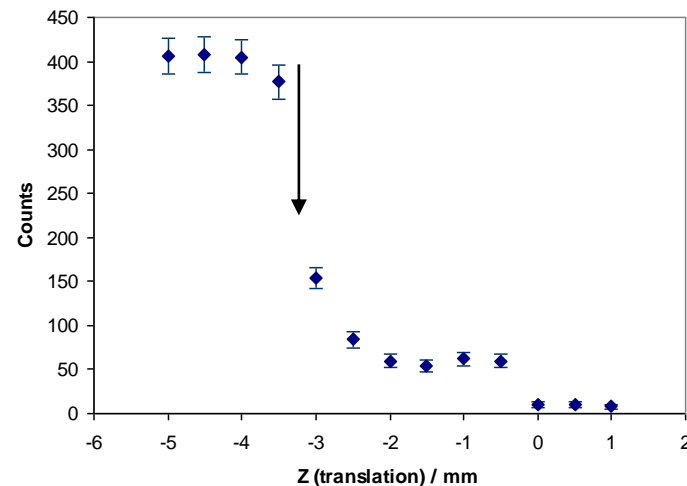
Choose fine slits to give collimated beam



Scan z

Look at intensity on detector

Identify $z = -3.2$ (~ 230 cts) as position
interface intersects direct beam

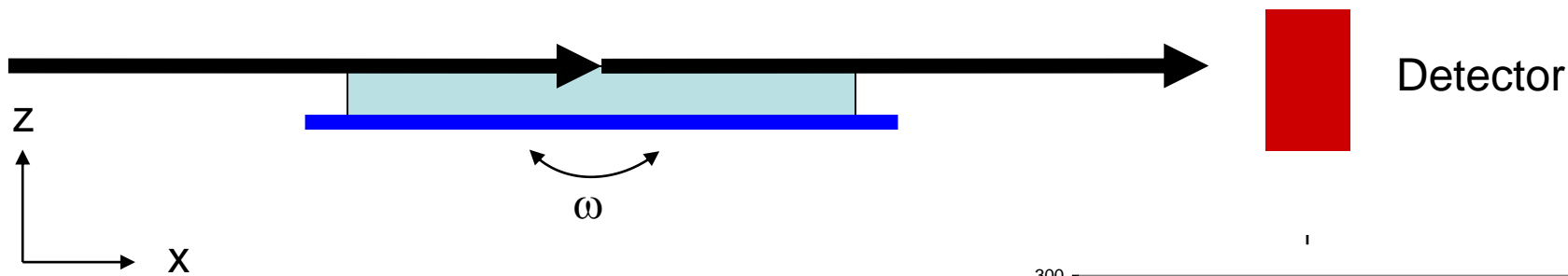




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Aligning a Sample

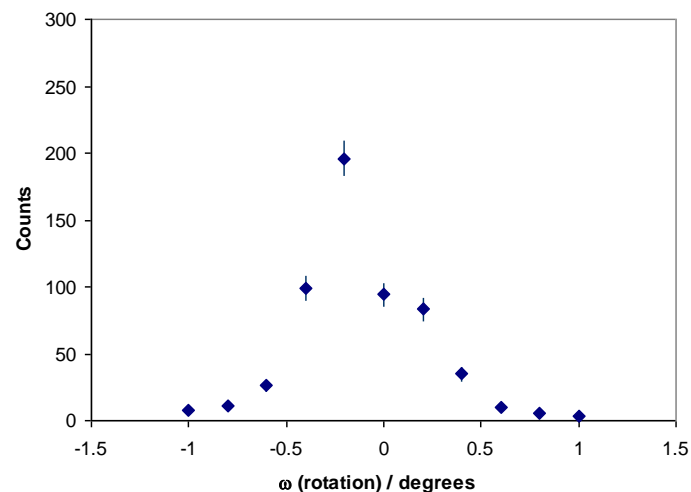
Move z to approximate sample in beam position



Scan ω

Look at intensity on detector

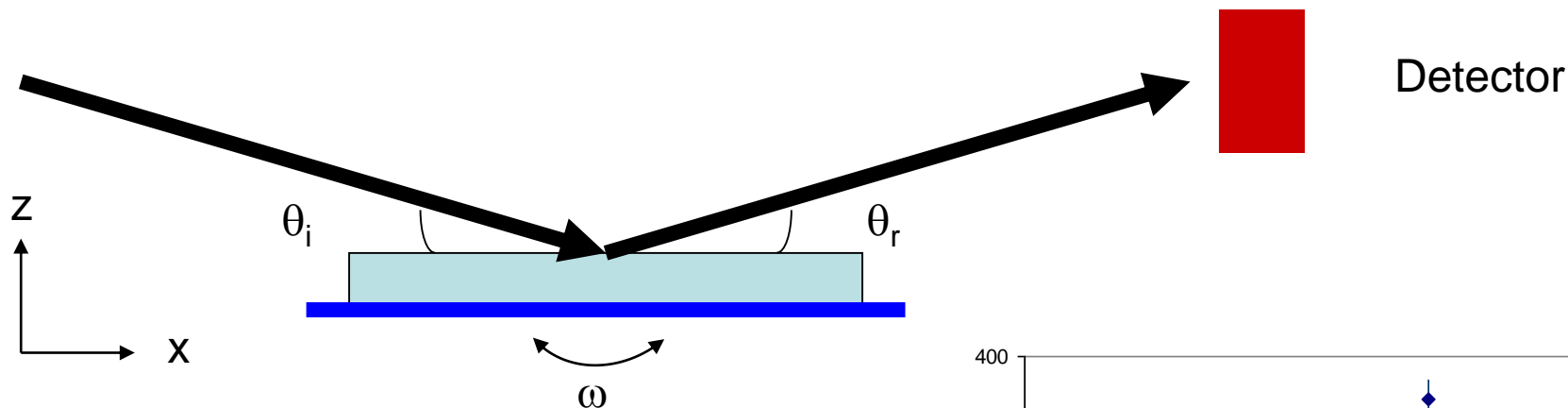
Identify $\omega = -0.22$ (~190 cts) as approximate sample offset angle





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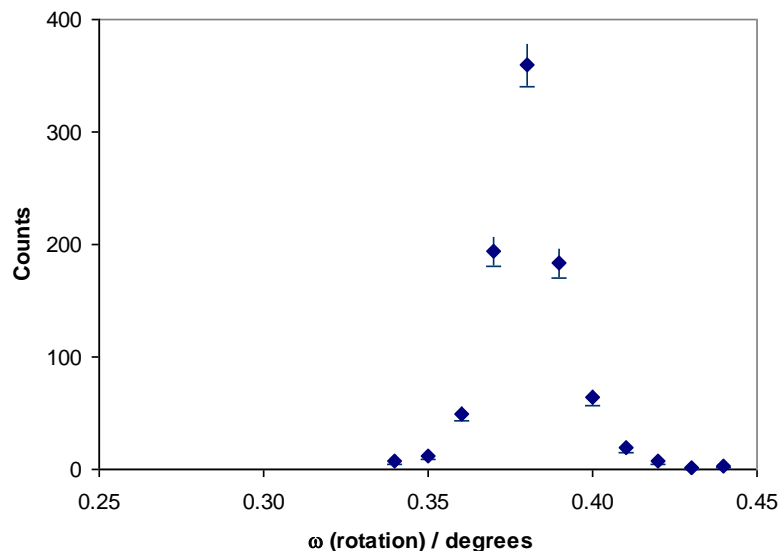
Aligning a Sample



Use approximate ω and z offset from alignment on direct beam

Set detector to small angle of reflection (e.g. 0.5°) and align more precisely.

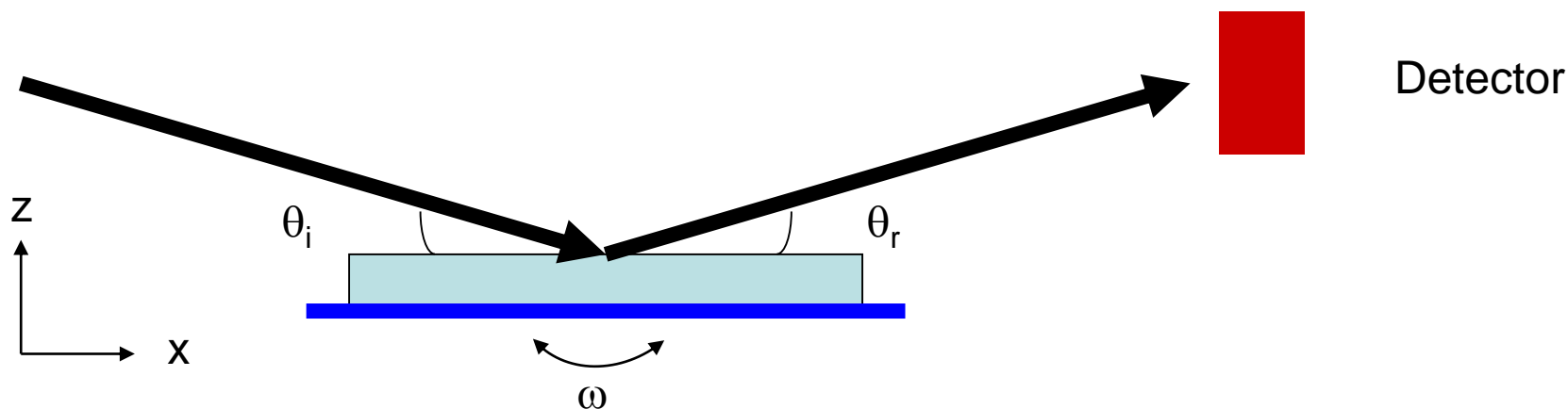
Scan ω and look for peak. Position is 0.378° and so offset is -0.122° .





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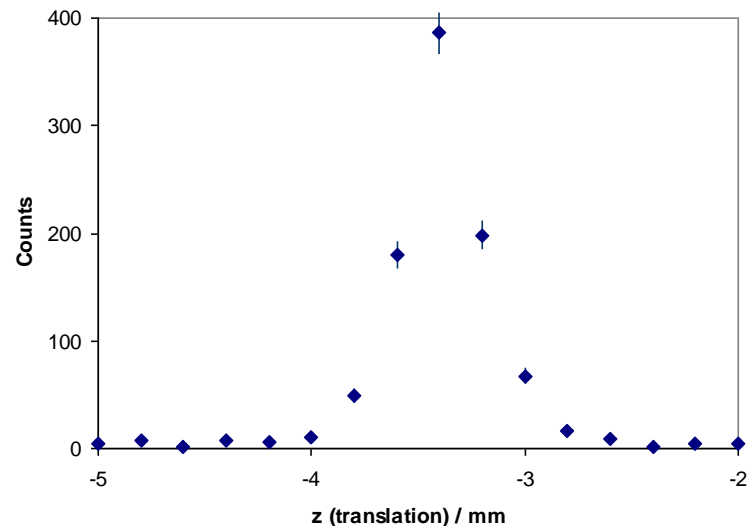
Aligning a Sample



Use new ω offset and z offset from alignment on direct beam

Check translation (z) offset in reflection mode.

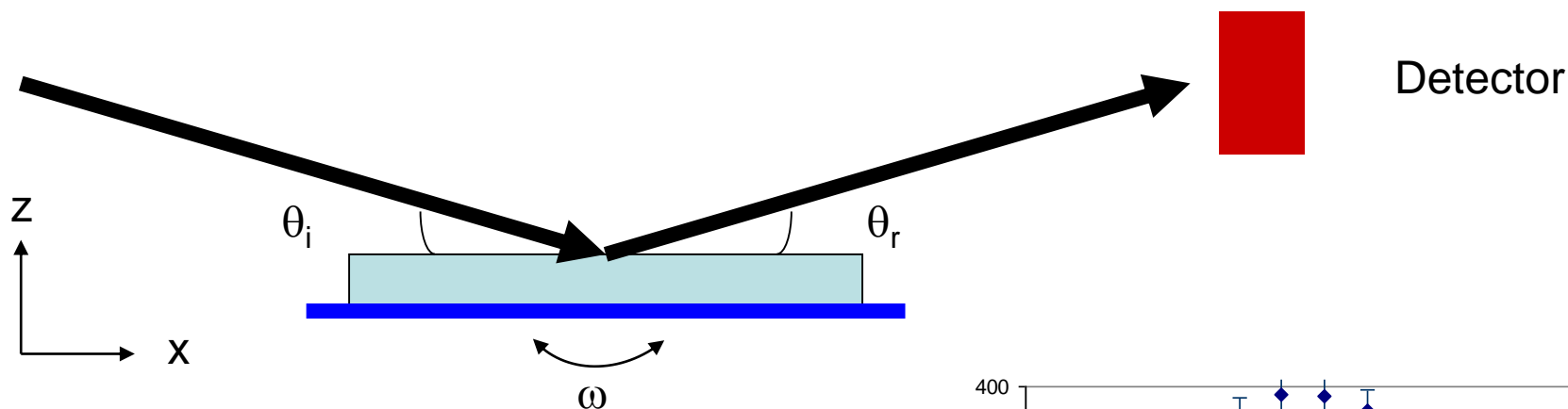
Scan z and look for peak. Position is -3.38 mm.





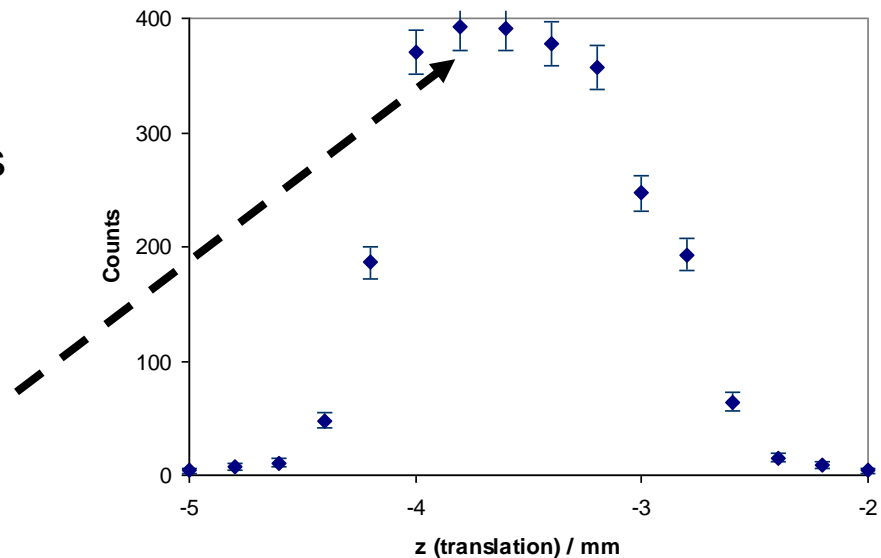
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Comments on Alignment



Angular (ω) width can depend on flatness of sample as well as resolution from slits and wavelength spread

If sample is very under-illuminated, translation (z) scan will have a flat top

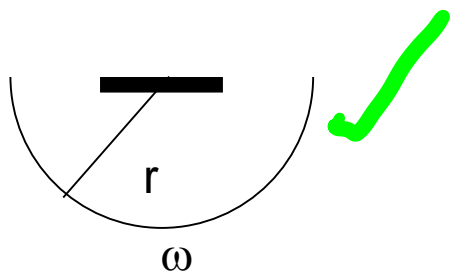




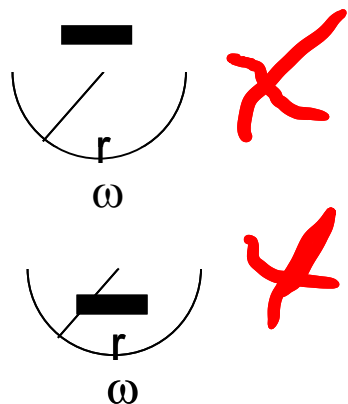
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Summary - Mounting and Alignment

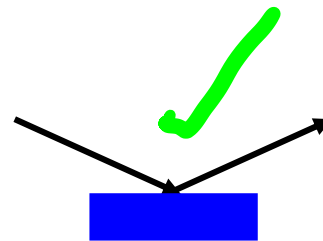
(a)



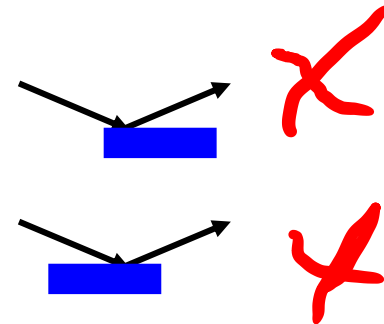
Eucentric mount



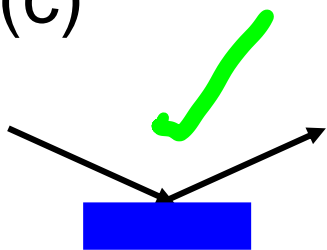
(b)



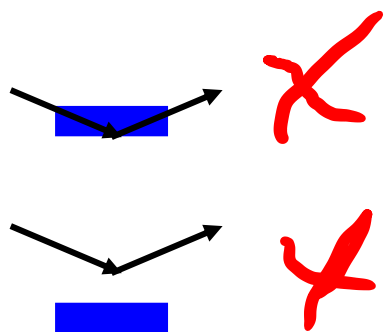
Centred on beam



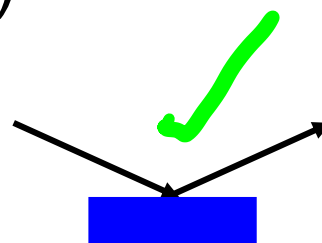
(c)



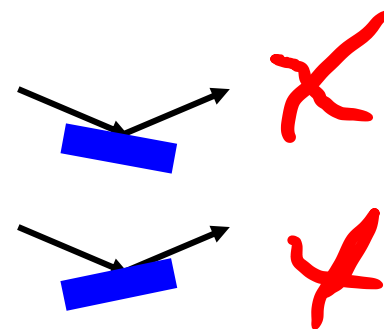
z position correct



(d)



ω offset correct





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Comments on Alignment

Using the results of alignment scans needs **offsets** or new zero positions to be set on the instrument. **Warning:** there is no general convention of signs on different instruments

Linear **thermal expansion** can be $\sim 2 \times 10^{-5}$ K⁻¹. 4 cm of aluminium changed by 50 C gives a shift of 0.04 mm.

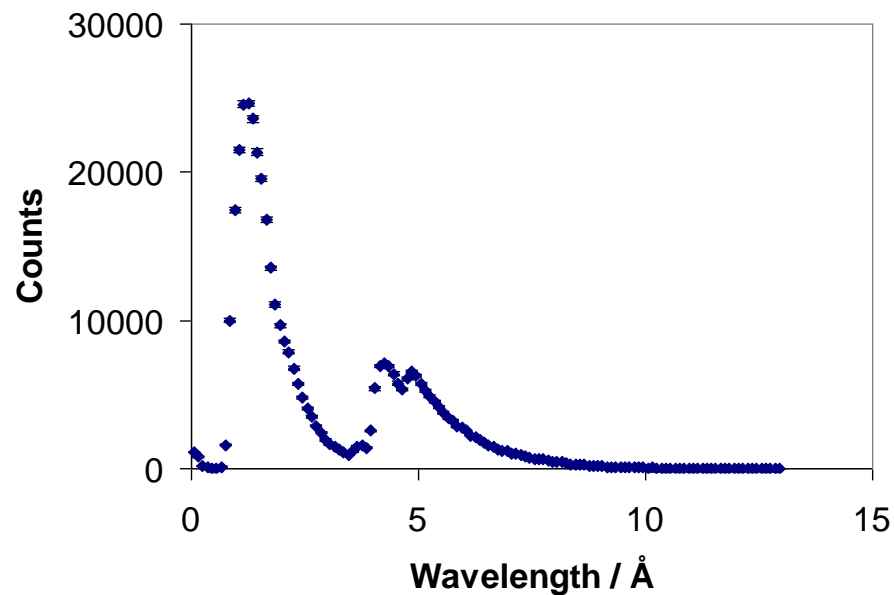


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Calibrations

Scan angle, measure
different λ or a combination
of λ and angle

Measure direct beam
(through sample
environment if needed)



Incident beam spectrum, LARMOR



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Samples

Low incident angle requires large uniform surface area. Footprint $\sim s / \tan \theta$.

Areas often several cm^2 .

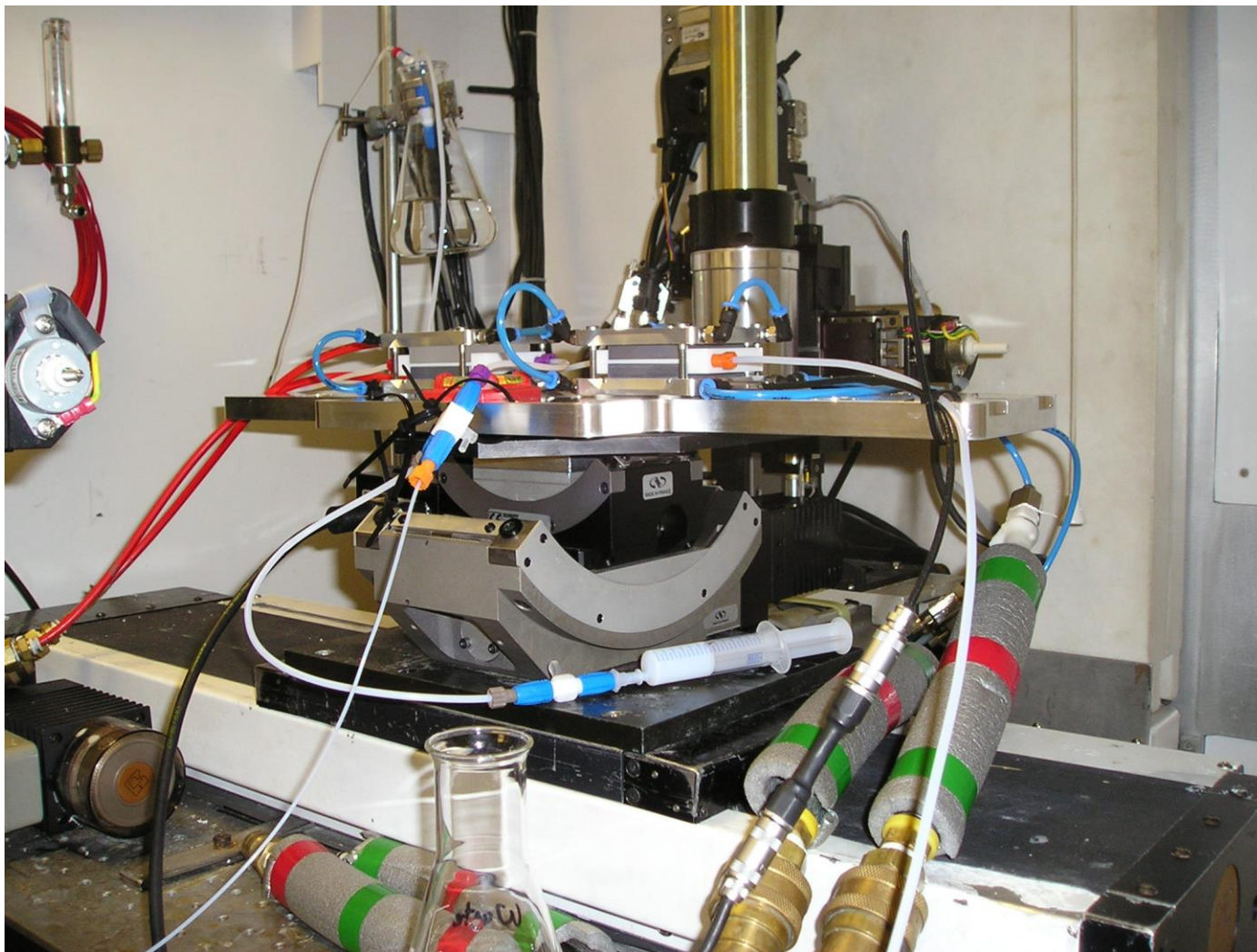
Smooth surface. 10 \AA roughness will reduce the reflectivity at $q=0.1 \text{ \AA}^{-1}$ by 2.7. 15 \AA reduces reflectivity by a factor of 10.

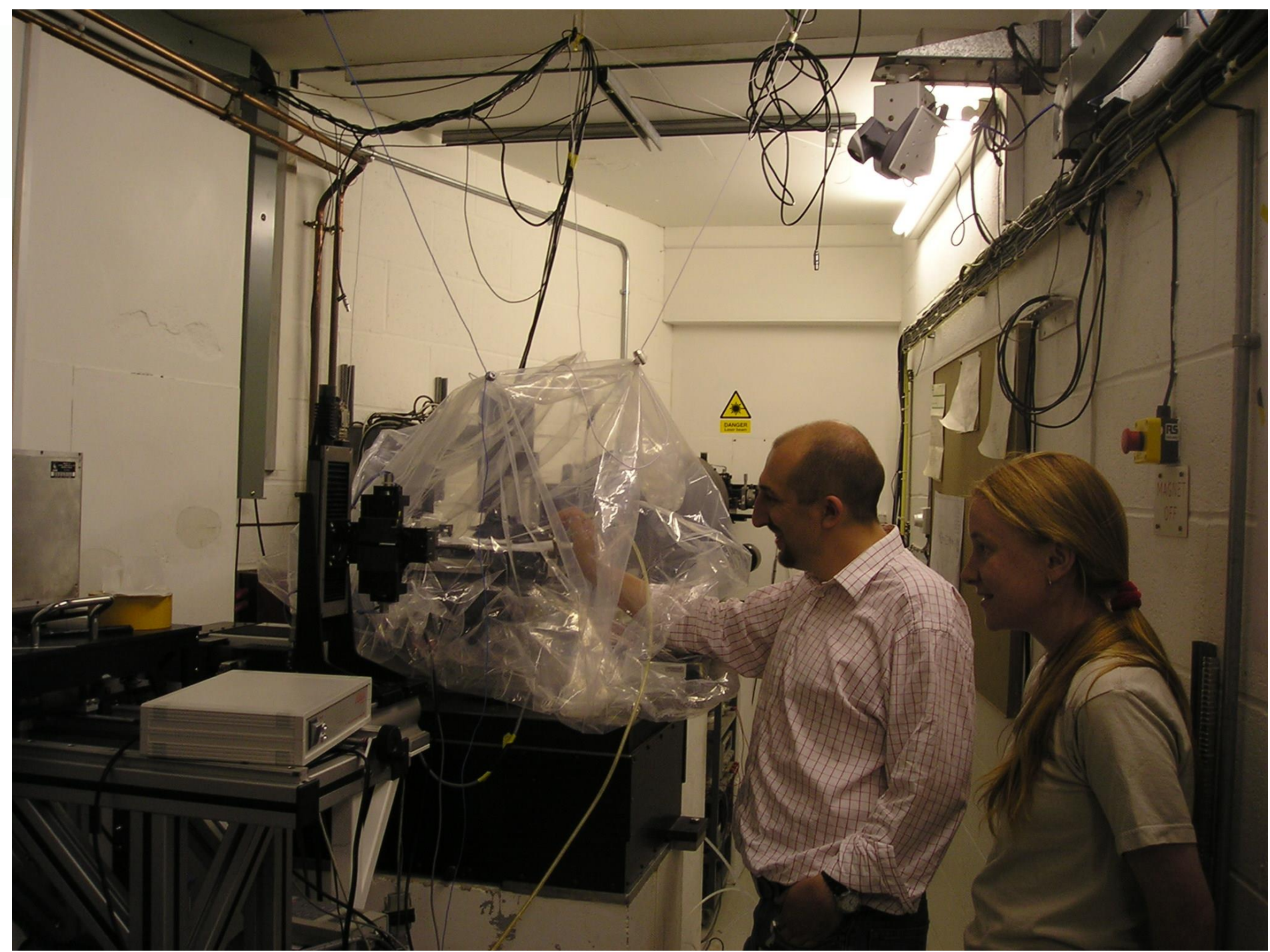
Liquids will have surface oscillations (capillary waves). Need to avoid other, induced waves.



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Sample Cell



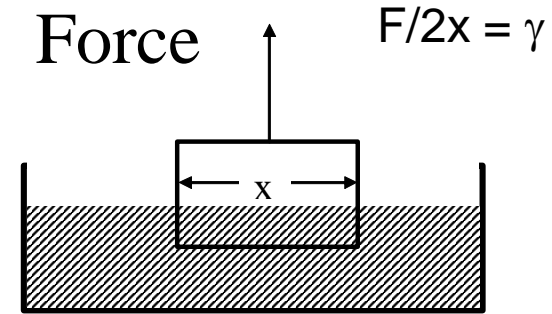
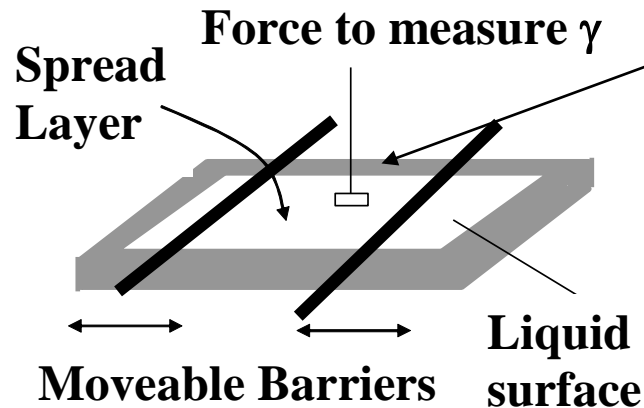




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Chemistry on a Liquid Surface

Langmuir Trough



In place of a drop use,
a uniform flat surface



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What is measured?

Reflected signal may have a large background

For hydrogenous substrate $\sim 5 \times 10^{-6}$
incident beam

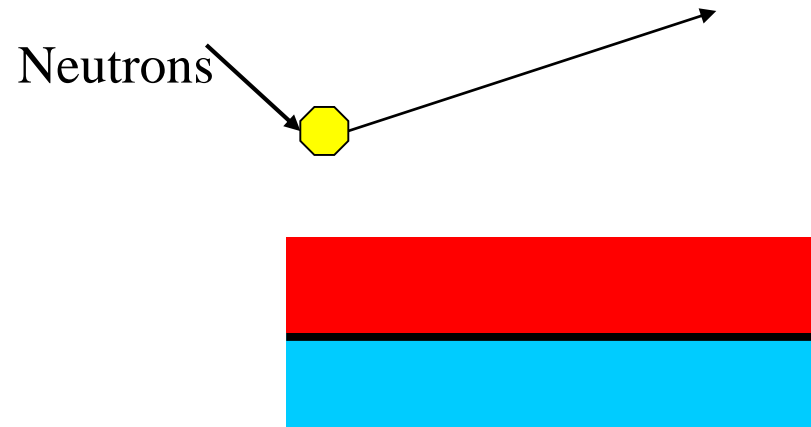
Attenuation by reduced transmission
(caused by scattering or absorption) may
be significant



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Fate of a Neutron at an Interface

- Reflected
- Scattered/Diffracted from surface
- Absorbed
- Scattered from bulk (either side of surface)
- Other accidents





What does background look like?

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X-ray scattering – glass

Sinha et al., *Phys. Rev. B.*

38, 2297, 1988.

Neutron scattering from
 D_2O and from null
reflecting water

Rennie et al.,

Macromolecules **22**, 3466-
3475 (1989).

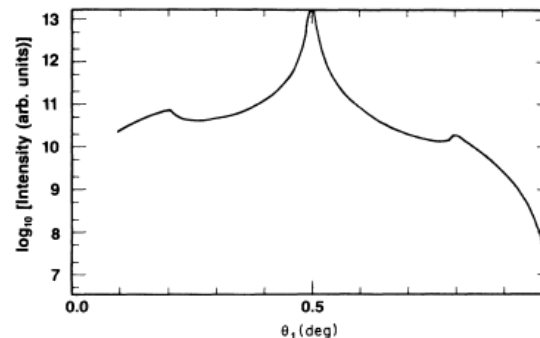
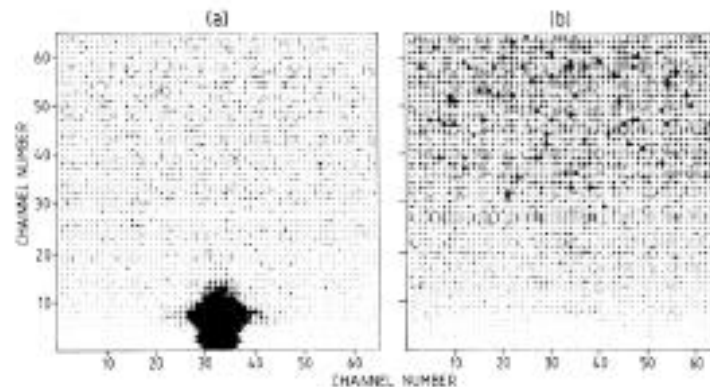


FIG. 6. Calculation of diffuse scattering in the distorted-wave Born approximation for rocking curve where θ_1 and θ_2 are varied such that 2θ is fixed at 1° . The asymmetry is due to the area of the illuminated surface decreasing as θ_1 is increased. The q_y direction has been integrated over. Parameters are $\sigma = 7 \text{ \AA}$, $h = 0.2$, $\xi = 7000 \text{ \AA}$, and the optical constants for Pyrex are given in Sec. V.





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Contrast Matching

$$\text{H}_2\text{O} \quad \rho = -0.56 \times 10^{-6} \text{ \AA}^{-2}$$

$$\text{D}_2\text{O} \quad \rho = +6.35 \times 10^{-6} \text{ \AA}^{-2}$$

$$y \times 6.35 + (1-y) \times (-0.56) = 0$$

$$6.91 y = 0.56 \quad \text{or} \quad y = 0.56 / 6.91 = 0.081$$

i.e. 8% by volume of D_2O in H_2O has $n = 1$



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Comments on Calculations

Programs that lose data

It is common to use logarithmic scales but background subtraction can give negative data points. $R Q^4$ is useful.

Experimental issues

Resolution – often needs to be included

Illumination

Small samples are often not able to reflect all the beam and a geometrical correction is applied.

Absolute reflectivity

Data is constrained if it is on an absolute scale



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What has not (yet) been covered?

Ellipsometry and X-rays

Needs more calculations for s and p waves

How to write a minimisation routine?

How to install your favourite program?

Specific examples of real samples etc.



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Do's and Don'ts

- Do not bend samples – care with mounts
- Use anti-vibration mounts for liquids – air borne noise causes vibrations
- Capillary waves cause scattering



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Questions?



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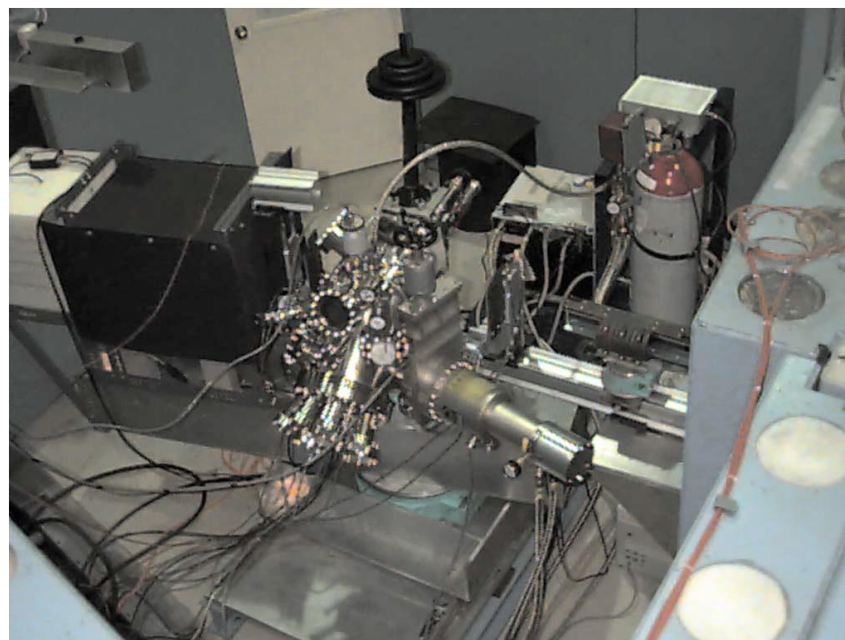
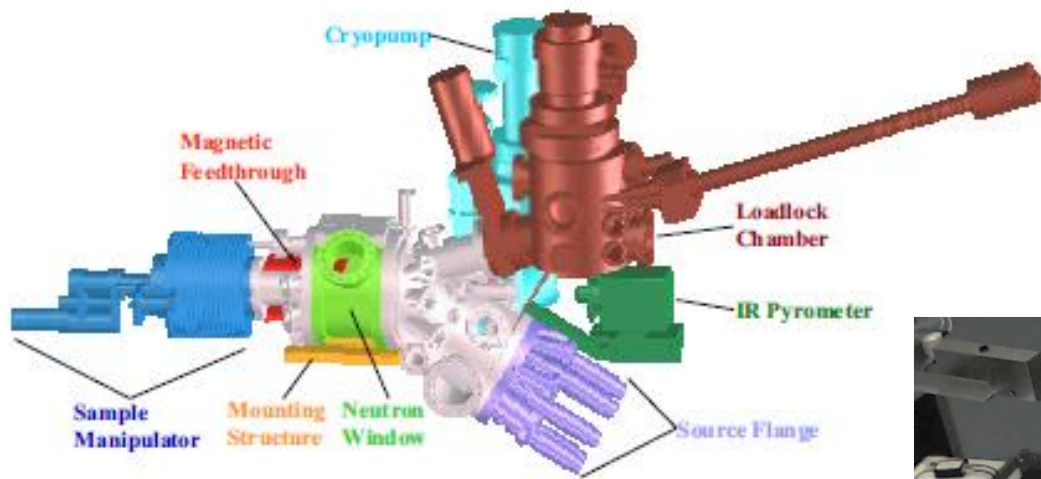
Sample Environments

- Choice is very large
- Build for your own experiment



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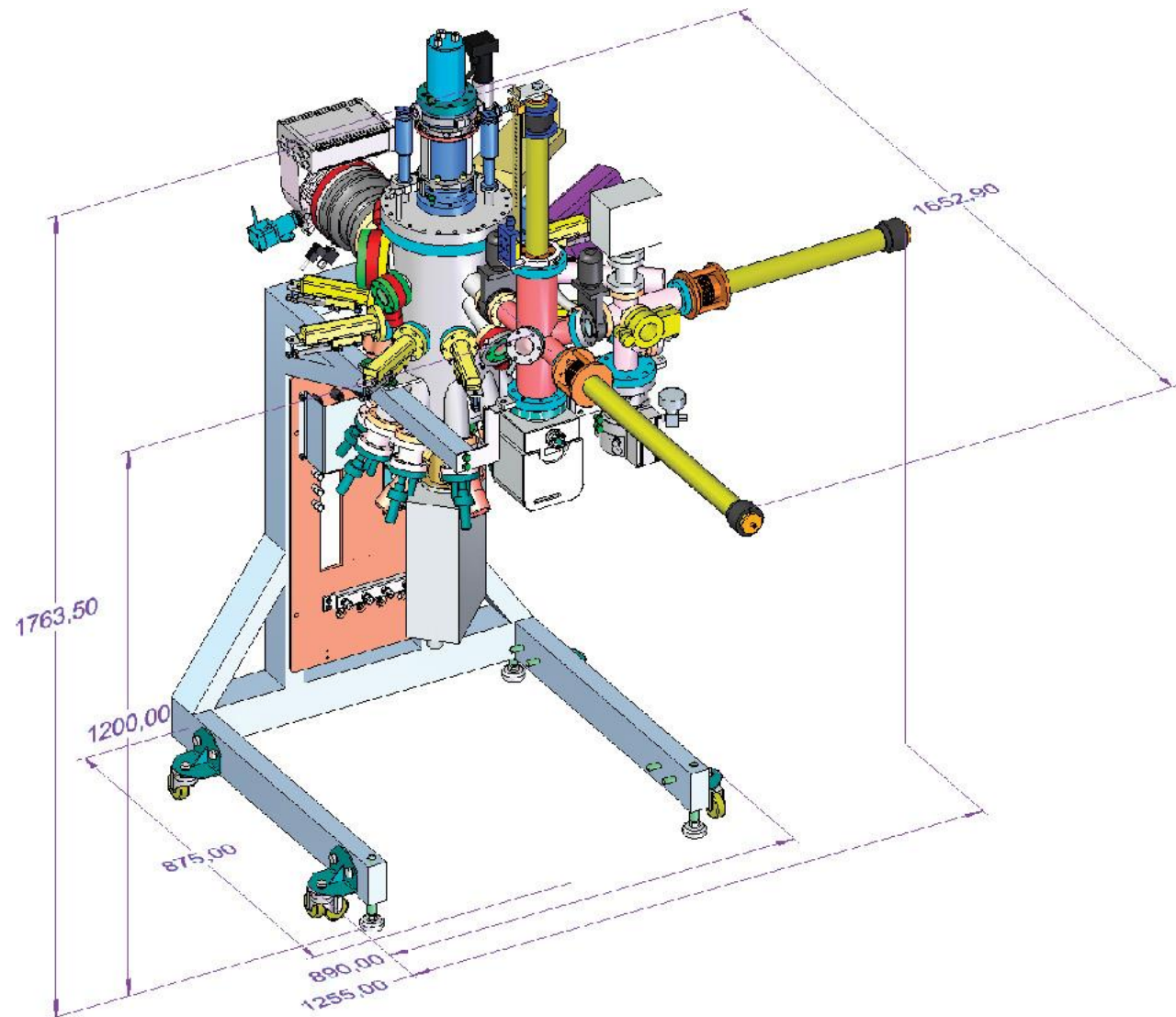
Thin Film Growth



J. A. Dura, J. LaRock 'A molecular beam epitaxy facility for in situ neutron scattering' *Rev. Sci. Instrum.* **80**, (2009), 073906.



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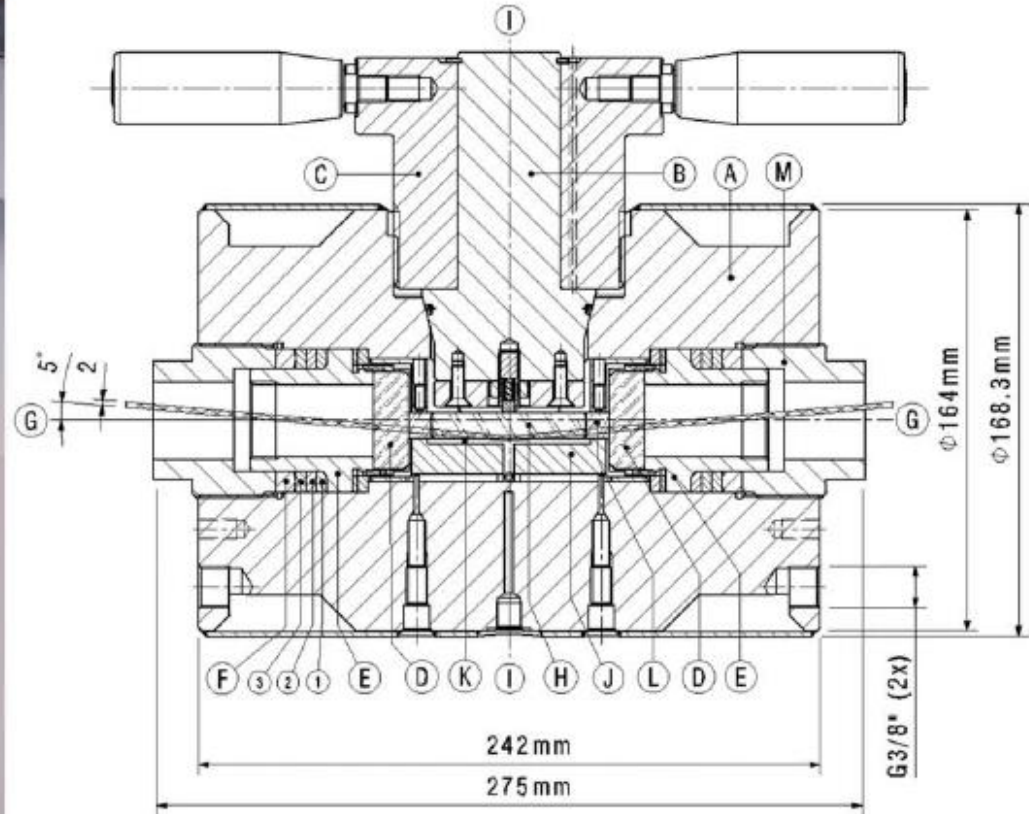


A. A. Baker, W. Braun, G. Gassler, S. Rembold, A. Fischer, T. Hesjedal 'An ultra-compact, high-throughput molecular beam epitaxy growth system' *Review of Scientific Instruments* **86**, (2015), 043901.



High Pressure

U

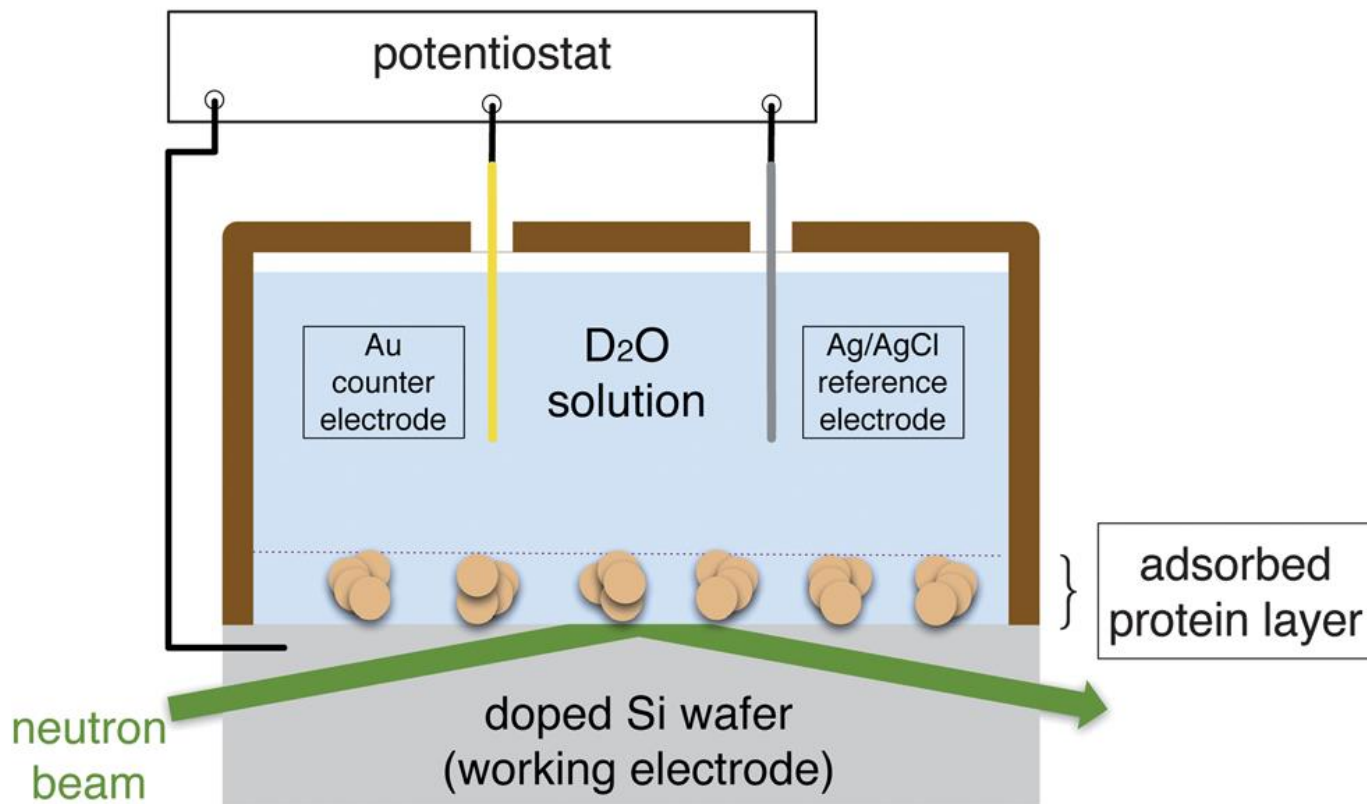


Martin Kreuzer, Thomas Kaltoven, Roland Steitz, Beat H. Zehnder, Reiner Dahint 'Pressure cell for investigations of solid-liquid interfaces by neutron reflectivity' *Rev. Sci. Instrum.* **82**, (2011), 023902.



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Electric Potential and Electric Currents

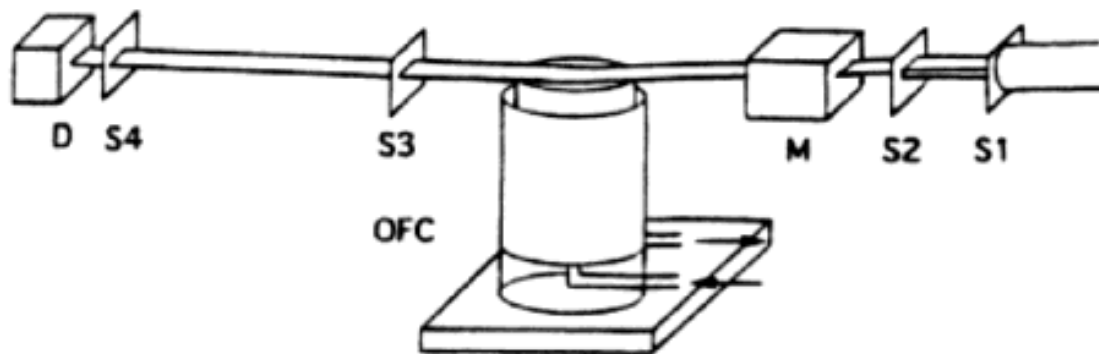


Alexandros Koutsioubas, Didier Lairez, Gilbert Zalczner, Fabrice Cousin 'Slow and remanent electric polarization of adsorbed BSA layer evidenced by neutron reflection' *Soft Matter*, **8**, (2012), 2638-2643.

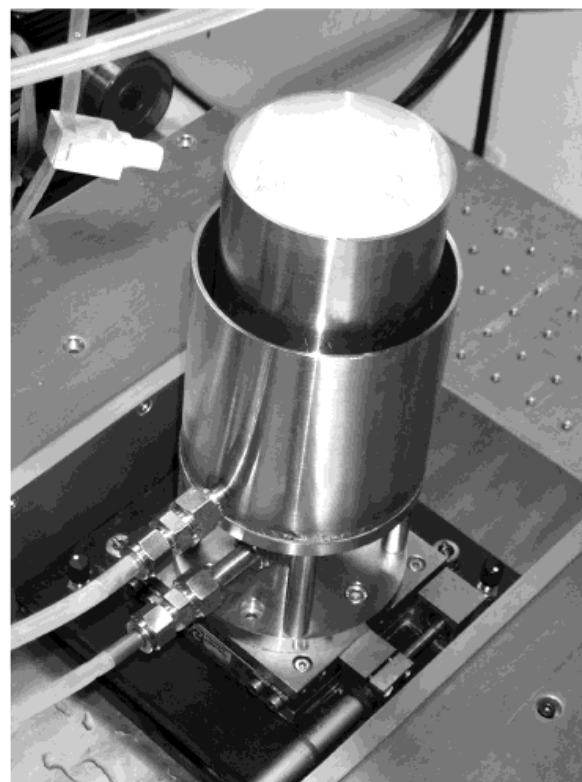


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Continuously Generated Fresh Liquid Surface



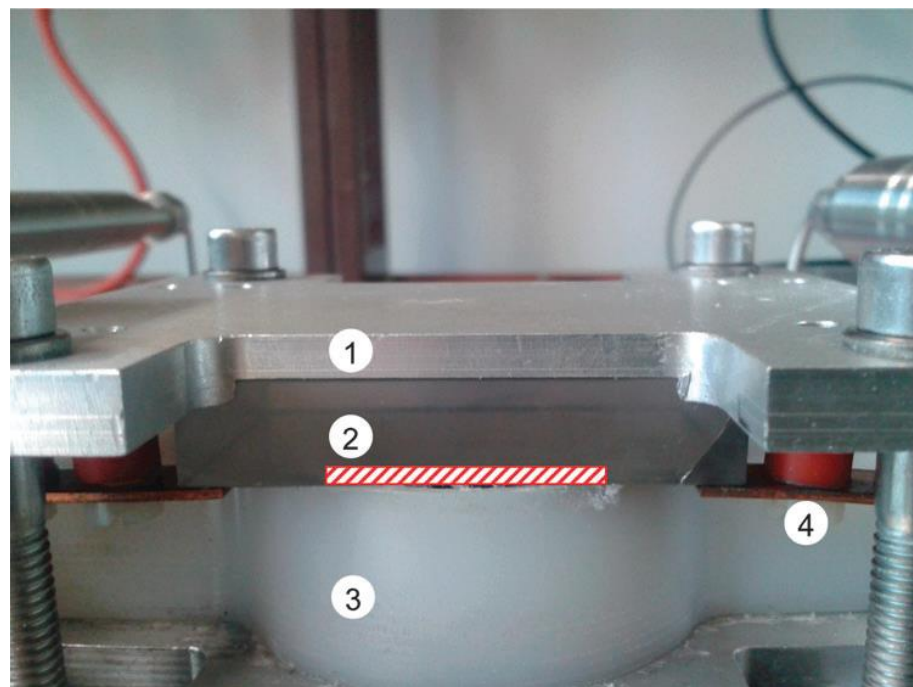
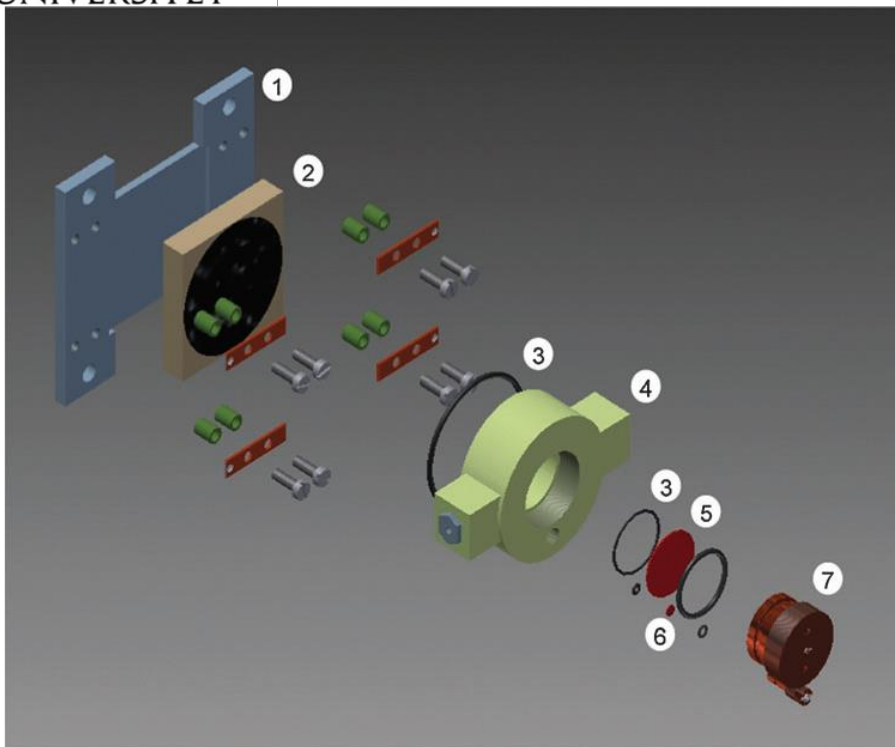
Julian Eastoe, Alex Rankin, Ray Wat, Colin D. Bain,
Dmitrii Styrkas, Jeff Penfold 'Dynamic Surface
Excesses of Fluorocarbon Surfactants' *Langmuir*,
19, (2003), 7734-7739.





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Battery Electrodes



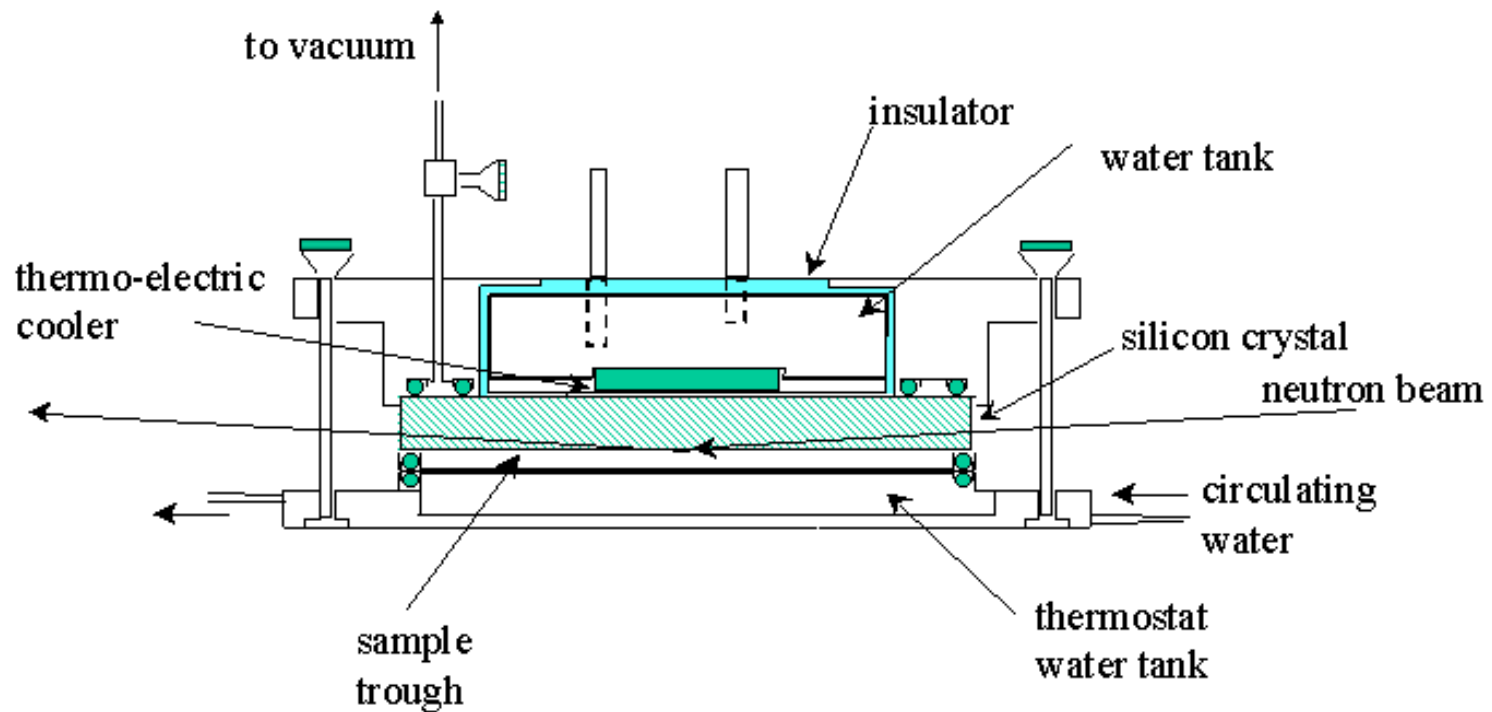
B. Jerliu, L. Dörrer, E. Hüger, G. Borchardt, R. Steitz, U. Geckle, V. Oberst, M. Bruns, O. Schneider, H. Schmidt 'Neutron reflectometry studies on the lithiation of amorphous silicon electrodes in lithium-ion batteries' *Phys. Chem. Chem. Phys.*, **15**, (2013), 7777-7784.



Liquid / Liquid Interfaces

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A. Zarbakhsh, J. Bowers, J. R. P. Webster, 'A new approach for measuring neutron reflection from a liquid/liquid interface' *Meas. Sci. Technol.* **10**, (1999), 738-743.



Other Ideas and Possibilities

http://www.reflectometry.net/reflect_bib.htm#Sample_environment

Sample Environment for Reflection

Seq. No.	Reference	Digital Source - DOI	Year	Technique
61	F. A. Adlmann, P. Gutfreund, J. F. Ankner, J. F. Browning, A. Parizzi, B. Vacaliuc, C. E. Halbert, J. P. Rich, A. J. C. Dennison, M. Wolff 'Towards neutron scattering experiments with submillisecond time resolution' <i>J. Appl. Cryst.</i> 48 , (2015), 220-226.	http://dx.doi.org/10.1107/S1600576714027848	2015	Oscillatory Shear
42	Anna Angus-Smyth, Richard A. Campbell, Colin D. Bain 'Dynamic Adsorption of Weakly Interacting Polymer/Surfactant Mixtures at the Air/Water Interface' <i>Langmuir</i> , 28 , (2012), 12479-12492.	http://dx.doi.org/10.1021/la301297s	2012	Overflowing Cylinder
15	I. F. Bailey 'A review of sample environments in neutron scattering' <i>Z. Kristallogr.</i> 218 , (2003), 84-95.	http://dx.doi.org/10.1524/zkri.218.2.84.20671	2003	Review
5	Shenda M. Baker, Gregory Smith, Roger Pynn, Paul Butler, John Hayter, William Hamilton, Lee Magid 'Shear cell for the study of liquid-solid interfaces by neutron scattering' <i>Rev. Sci. Instrum.</i> 65 , (1994), 412-416.	http://dx.doi.org/10.1063/1.1145148	1994	Shear
59	A. A. Baker, W. Braun, G. Gassler, S. Rembold, A. Fischer, T. Hesjedal 'An ultra-compact, high-throughput molecular beam epitaxy growth system' <i>Review of Scientific Instruments</i> 86 , (2015), 043901.	http://dx.doi.org/10.1063/1.4917009	2015	MBE
4	T. M. Bayerl, R. K. Thomas, A. R. Rennie, J. Penfold, E. Sackmann, 'Specular reflection of neutrons at phospholipid monolayers: changes of monolayer structure and head group hydration at the transition from the expanded to the condensed phase state', <i>Biophysical Journal</i> 57 , (1990), 1095-1098.	http://dx.doi.org/10.1016/S0006-3495(90)82628-X	1990	Langmuir trough
62	N. Booth, G. Davidson, P. Imperia, S. Lee, B. Stuart, P. Thomas, K. Komatsu, R. Yamane, S. W. Prescott, H. E. Maynard-Casely, A. Nelson, K. C. Rule 'Three impossible things before lunch – the task of a sample environment specialist' <i>Journal of Neutron Research</i> 19 , (2017), 49-56.	http://dx.doi.org/10.3233/JNR-170041	2017	Conductivity, in-line
13	James Bowers, Ali Zarbakhsh, John R. P. Webster, Lian R. Hutchings, Randal W. Richards 'Neutron Reflectivity Studies at Liquid-Liquid	http://dx.doi.org/10.1021/la001119o	2001	Liquid/liquid interface



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Questions?

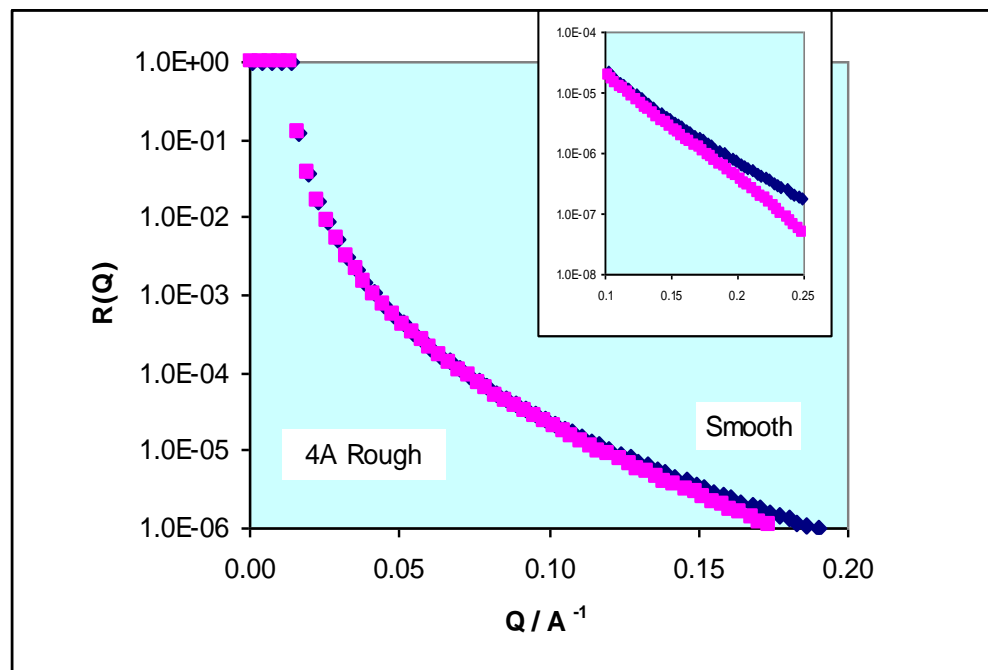


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Roughness

Reflectivity from rough surfaces is decreased.

‘Gaussian’ roughness’
– intensity decreases
by $\exp(-Q^2\xi^2/2)$ for
scattering vector, Q
and amplitude of
roughness, ξ .





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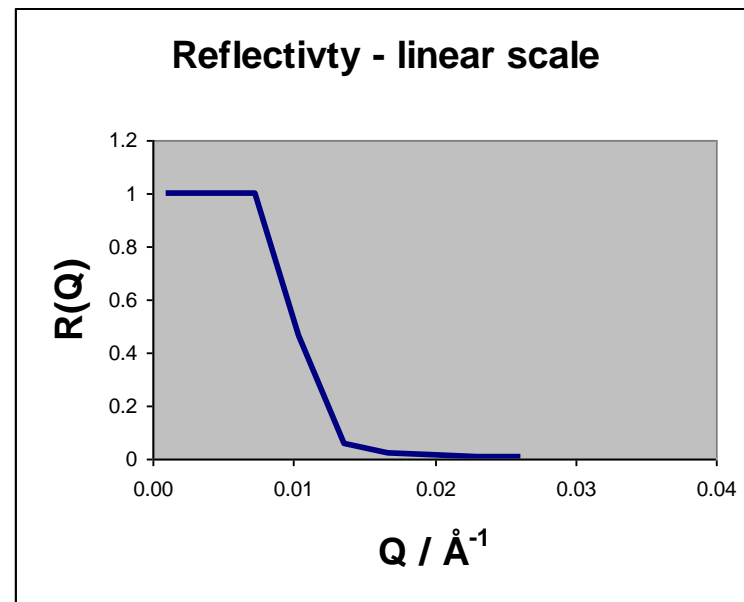
Critical Angle and Below (critical wavelength and above)

Density difference between two bulk phases determines the critical momentum transfer/angle, Q_c or θ_c

Any variation in intensity below critical angle is probably telling you about the experiment rather than the interface

$R = 1$ for $\theta < \theta_c$ is often used as a calibrant

Total reflection below critical angle θ
 $\cos \theta = n_2/n_1$





Intensity of Reflected Signal

- Waves interfere constructively for
$$2 d \sin \theta = \lambda, 2\lambda, 3\lambda \dots$$
- Measured reflectivity will depend on angle and wavelength. Add wave amplitudes with allowance for phase and calculate intensity as square of amplitude.
- Total reflection for angles less than critical angle, $\theta_c = \arccos(n_1/n_2)$



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Fresnel Formula

Reflection from an interface between two media with $\Delta\rho = \rho_1 - \rho_2$ is for $Q \gg Q_c$:

$$R(Q) = 16 \pi^2 (\Delta\rho)^2 / Q^4$$

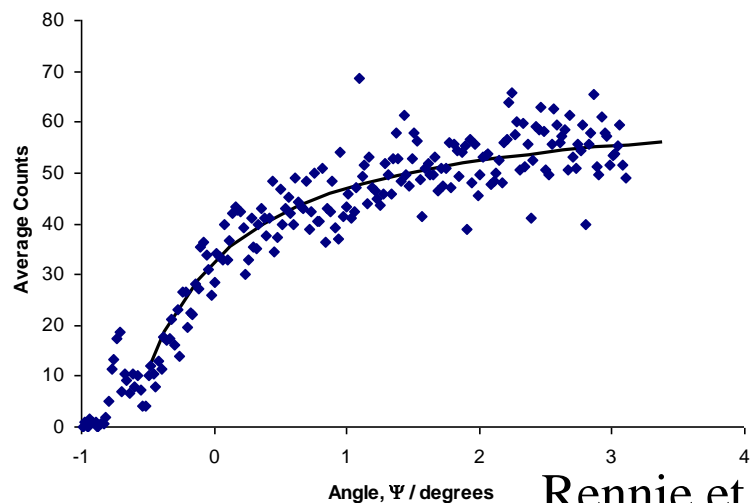
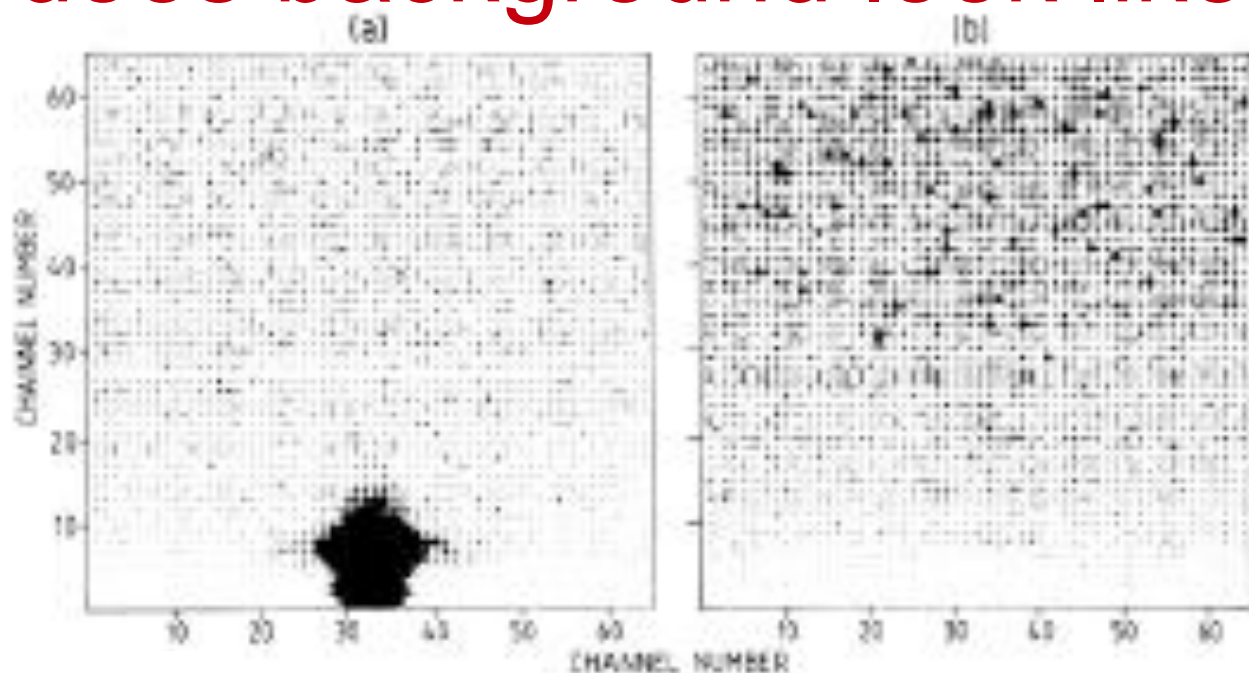
Note

This does not depend on sign of $\Delta\rho$.



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What does background look like?



Scattering from D_2O
and from null reflecting water
(8% D_2O)

Rennie et al., *Macromolecules* **22**, (1989), 3466-3475.